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Constraining the fireball scenario of GRB afterglows with GROND and multi-wavelength data.

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We performed an analysis of the multi-epoch broad-band observations of the GRB 121024A afterglow covering the full range from radio to X-rays. From the temporal and spectral evolution of the afterglow we set constraints on the micro-physical and dynamical parameters describing the physics of the afterglow. We tested both a jet break scenario and an energy injection model for the interpretation of the break in the light curve. The jet break model has been suggested by Wiersema et al., based on the linear and circular polarisation detections for this burst. The final values for the micro-physical and dynamical parameters in this model are physically plausible. However, it requires a hard distribution for the shock-accelerated electron energies, an extremely low value for the cooling break frequency, a non-spreading jet and an extreme prompt emission efficiency. The energy injection model avoids the unusual requirements of the jet break model but it gives some atypical values for the micro-physical and dynamical parameters.
1. Introduction

In the standard GRB afterglow model [1], the observed emission is associated with synchrotron radiation from accelerated electrons. The observed spectrum has 3 characteristic breaks: the cooling frequency $\nu_c$, the synchrotron injection frequency $\nu_m$ and the self-absorption frequency $\nu_{sa}$. Each one yields specific and correlated constraints on the physical processes in the relativistic outflow. These processes are described by 3 micro-physical parameters (i.e. fraction of energy in the electrons $\epsilon_e$ and in the magnetic field $\epsilon_B$, power-law index of the non-thermal electron population $p$) and 3 dynamical parameters (i.e. isotropic equivalent energy $E_{iso}$, circumburst density profile $n = Ar^{-k}$ with $A$ a constant and $r$ the radius, half-opening angle $\theta_0$). Here, we present the analysis of the simultaneous multi-wavelength observations of GRB 121024A. It was followed up by different instruments in wavelengths from radio to X-rays during several days, and has a redshift $z = 2.30$ measured by the X-shooter spectrograph at the Very Large Telescope (VLT) [2]. From these simultaneous broad-band observations, we derive constraints on the micro-physical and dynamical parameters of the GRB afterglow. This burst is well known as the first one for which linear and circular optical polarisations have been observed that were claimed to cover a jet break [3]. On the other hand a prolonged energy injection model avoids the main issues of the jet break scenario but gives atypical values for some of the derived parameters. Here, we present both scenarios.

2. Observations and data reduction

**Swift X-ray Telescope (XRT):** On October 24th 2012 at 02:56:12 UT the Swift Burst Alert Telescope (BAT) triggered and located GRB 121024A [4]. XRT slewed immediately to the burst and the observations started 93 sec after the trigger. The Swift/XRT light curve and spectra data were obtained from the XRT repository [5]. The measured energy in the prompt emission is $E_{\gamma,iso} = 8.4^{+2.6}_{-2.2} \times 10^{52}$ erg [6].

The **Gamma-Ray burst Optical Near-infrared Detector - GROND** [7] started observations of the GRB field 2.96 hours after the Swift trigger [8] and continued for the next 3.8 hours. The afterglow was detected at RA(J2000) = 04:41:53.30 and Dec(J2000) = -12:17:26.5 with an uncertainty of 0.4 in each coordinate in all the 7 bands $g'r'i'z'JHK_s$. Imaging of the field of GRB 121024A continued on the 2nd, 3rd, 4th, 16th and 17th night after the burst. The optical/NIR data was reduced using standard IRAF tasks [9]. The optical magnitudes were calibrated against secondary stars in the GRB field. The NIR magnitudes were calibrated against the Two Micron Sky Survey -2MASS- [10] catalogue stars in the field of the GRB.

The **Large APEX Bolometer Camera LABOCA** [11] was triggered on October 24th 2012. Two observations at a frequency of 345 GHz with a band width of 60 GHz were performed. The first one started 19.8 ks after the GRB and the second was at a mid-time of 109.0 ks after the trigger. During both days, observations were taken in mapping and in on-off mode. There was no detection in either of both nights. The upper limits are 3.6 mJy beam$^{-1}$ and 10.4 mJy beam$^{-1}$ for the 1st and 2nd nights, respectively.

**Millimeter and radio observations:** The following values published in the literature are used in the analysis of the complete broad-band energy spectrum: The Combined Array for Research in
Millimeter-wave Astronomy (CARMA) started observations of the field of GRB 121024A ∼ 120.9 ks after the BAT trigger at a mean frequency of ∼ 85 GHz (3mm) [12]. A mm counterpart was detected with a flux of 1.0 ± 0.3 mJy. The Very Large Array (VLA) started observations of the field of GRB 121024A ∼ 109.0 ks after the trigger. They were performed using the K, C and X bands. A radio counterpart with flux of 0.10 ± 0.03 mJy was detected at frequency of 22 GHz [13].

3. Results

To study the temporal evolution, we first perform a combined fit using XRT and GROND data using a smoothly broken power-law with host contribution. The best fit reveals an achromatic break with slope going from $\alpha_{\text{pre}} = -0.85 \pm 0.04$ to $\alpha_{\text{post}} = -1.47 \pm 0.04$, smoothness $s_m = 1.7 \pm 0.3$ and break time $t_b = 49.8 \pm 5.1$ ks.

To test for spectral evolution, we first use X-ray and optical/NIR measurements at 6 different epochs (see Fig. 2a). The best fitting profile is a power-law with $\beta = 0.86 \pm 0.02$, showing no evidence for spectral evolution. There are two possible scenarios that can explain the observations. On one hand, we have $\nu_{\text{NIR}} > \nu_{\text{c}}$ with $p = 1.73 \pm 0.03$. This is a flat electron spectrum with the break in the light curve associated with a jet break. On the other hand, we obtain $\nu_{\text{c}} > \nu_{\text{x-rays}}$ with $p = 2.73 \pm 0.03$. In this case, the break in the light curve is the end of the energy injection phase in a stellar-wind type environment.

Now, we proceed to used the broad-band observations (Fig. 2b). For the SED at $t=19.8$ ks the APEX upper limit implies a break between APEX and NIR bands. This break is associated with $\nu_{\text{c}}$ in the jet break scenario and with $\nu_{\text{m}}$ in the energy injection scenario. For the SED at $t=109.0$ ks, the CARMA data point requires at least one break between this wavelength and the NIR bands and the EVLA data point implies a break between this wavelength and the CARMA wavelength. For the energy injection scenario, this break is associated with $\nu_{\text{sa}}$. We analyse this scenario for ISM, wind and a generic density profile with slope 1.1. For the jet break scenario, we have two breaks, associated with $\nu_{\text{sa}}$ and $\nu_{\text{m}}$. In this scenario, $\nu_{\text{m}}$ can be above or below $\nu_{\text{sa}}$. We present both spectral regimes. The results for the parameters are presented in Table 1, where GS [14] and
DC [15] two different jet break scenarios for \( p < 2 \). The models presented here are those that were not ruled out by the closure relations.

![Image](image_url)

**Figure 2:** Left: Spectral energy distribution for the 6 SED highlighted in Fig. 1. SEDs I - IV are from data before the observed break in the light curve. SEDs V - VI are from data taken after the break. The SEDs are scaled with arbitrary factor for clarity in the plot. Right: The red line corresponds to the SED at \( t = 19.8 \) ks. The green line is the SED at \( t = 109.0 \) ks. EVLA, CARMA, GROND and XRT data are included. We measured \( v_{\text{sa}} \) and \( v_{\text{in}} \) for the slow cooling regime when \( v_{e} > v_{\text{sa}} \).

<table>
<thead>
<tr>
<th>Description</th>
<th>( \varepsilon_{e} )</th>
<th>( \varepsilon_{B} )</th>
<th>( A_{\gamma}, n_{0} )</th>
<th>( \theta_{0} ) [rad]</th>
<th>( E_{\text{iso}}, 52 \text{[erg]} )</th>
<th>( \eta )</th>
</tr>
</thead>
<tbody>
<tr>
<td>GS scenario where ( p = 1.73 \pm 0.03 )</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
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<tr>
<td>( v_{\text{sa}} &lt; v_{\text{in}} )</td>
<td>( 1.9^{+1.1}_{-0.8} \times 10^{-2} )</td>
<td>( 2.2^{+0.6}_{-0.3} \times 10^{-2} )</td>
<td>( 1.4^{+1.2}_{-0.8} )</td>
<td>( &gt; 1.2 \times 10^{-1} )</td>
<td>( &gt; 3.4 )</td>
<td>( &lt; 71% )</td>
</tr>
<tr>
<td>( v_{\text{in}} &lt; v_{\text{sa}} )</td>
<td>( &lt; 8.8 \times 10^{-4} )</td>
<td>( &lt; 8.3 \times 10^{-2} )</td>
<td>( &gt; 0.8 )</td>
<td>( &gt; 3.4 )</td>
<td>( &lt; 71% )</td>
<td></td>
</tr>
</tbody>
</table>

**DC scenario where \( p = 1.73 \pm 0.03 \)**

<table>
<thead>
<tr>
<th>Description</th>
<th>( \varepsilon_{e} )</th>
<th>( \varepsilon_{B} )</th>
<th>( A_{\gamma}, n_{0} )</th>
<th>( \theta_{0} ) [rad]</th>
<th>( E_{\text{iso}}, 52 \text{[erg]} )</th>
<th>( \eta )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( v_{\text{sa}} &lt; v_{\text{in}} )</td>
<td>( 79.0^{+1.1}_{-6.3} \times 10^{-2} )</td>
<td>( 5.0^{+0.2}_{-0.1} \times 10^{-2} )</td>
<td>( 3.1^{+0.1}_{-0.0} )</td>
<td>( &gt; 1.0 \times 10^{-2} )</td>
<td>( &gt; 2.3 \times 10^{-1} )</td>
<td>( &lt; 97% )</td>
</tr>
<tr>
<td>( v_{\text{in}} &lt; v_{\text{sa}} )</td>
<td>( &lt; 10.6 \times 10^{-2} )</td>
<td>( &lt; 19.1 \times 10^{-2} )</td>
<td>( &gt; 0.3 )</td>
<td>( &gt; 3.4 )</td>
<td>( &lt; 71% )</td>
<td></td>
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</table>

**Energy injection scenario where \( p = 2.73 \pm 0.03 \)**

<table>
<thead>
<tr>
<th>Description</th>
<th>( \varepsilon_{e} )</th>
<th>( \varepsilon_{B} )</th>
<th>( A_{\gamma}, n_{0} )</th>
<th>( \theta_{0} ) [rad]</th>
<th>( E_{\text{iso}}, 52 \text{[erg]} )</th>
<th>( \eta )</th>
</tr>
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<tr>
<td>( v_{e} &gt; v_{\text{crit}}, k = 2 )</td>
<td>( &gt; 1.1 \times 10^{-6} )</td>
<td>( &lt; 1.2 \times 10^{-6} )</td>
<td>( &gt; 1.2 \times 10^{-6} )</td>
<td>( &gt; 1.0 )</td>
<td>( &gt; 2.4 )</td>
<td>( &lt; 77% )</td>
</tr>
<tr>
<td>( v_{e} &gt; v_{\text{crit}}, k = 1.1 )</td>
<td>( &gt; 7.4 \times 10^{-2} )</td>
<td>( &lt; 2.1 \times 10^{-9} )</td>
<td>( &gt; 4.3 \times 10^{-9} )</td>
<td>( &gt; 0.8 )</td>
<td>( &gt; 3.4 )</td>
<td>( &lt; 71% )</td>
</tr>
<tr>
<td>( v_{e} &gt; v_{\text{crit}}, k = 0 )</td>
<td>( &gt; 7.5 \times 10^{-2} )</td>
<td>( &lt; 2.2 \times 10^{-9} )</td>
<td>( &gt; 1.2 \times 10^{-9} )</td>
<td>( &gt; 0.8 )</td>
<td>( &gt; 3.7 )</td>
<td>( &lt; 69% )</td>
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</tbody>
</table>

**Table 1:** Physical parameters for GS, DC and Energy injections models. \( \varepsilon_{e} = \varepsilon_{e} \times \left( |p - 2| / (p - 1) \right) \). For \( k = 2 \) we report the density in terms of \( A_{\gamma} \), where \( A = M/4\pi r_{\text{sw}} = 5 \times 10^{11} \text{A}_{\gamma} \) g cm\(^{-3}\) [16]. \( \eta = E_{\text{iso}}^{2} / (E_{\text{iso}}^{2} + E_{\text{iso}}) \). [17, 18]

## 4. Discussion and conclusions

We analysed the multi-wavelength observations of the afterglow of GRB 121024A. The combined GROND and XRT data allow us to measure with high accuracy the spectral slope in this energy regime, and therefore the electron index \( p \). We describe our complete set of observations using two different models, where the break in the light curve marks either a jet break or the cessation of energy injection. The jet break model requires a hard electron spectrum with \( p < 2 \), a very low cooling break frequency \([19]\) and a non-spreading jet \([20, 21]\). The energy injection model does not require high efficiency values and is not in contradiction with Fermi acceleration predictions for \( p \), but gives atypical values for some of the micro-physical and dynamical parameters.
Although no clear preferred model describing the observations of this GRB afterglow emerges, we are able to rule out some of the possible models suggested by the closure relations [22]. First, we rule out the spectral regime where $\nu_{sa} < \nu_m$ for the jet break model. In the GS scenario [14], it is ruled out because the spectral evolution will never cross that regime in the slow cooling phase and in the DC scenario [15] it is ruled out because the time when $\nu_m$ crosses $\nu_{sa}$ is before the time of the studied SED ($t = 109$ ks). Second, we rule out the energy injection model for the wind density profile $k = 2$ because $\bar{\epsilon}_e$ has to be larger than one, which is not physically meaningful.

There are two possible models left describing the observations. First, energy injection model where we either take $k = 0$ (ISM) or fit $k = 1.1$ for the density profile. This model requires extremely high density values compared to theoretical expectations and previous measurements, and it also implies spherical outflow geometry. Second, the jet break model for the spectral regime where $\nu_m < \nu_{sa}$. This gives physically meaningful micro-physical and dynamical parameters, although it has some issues with the efficiency requirements, the position of the cooling break, and the hard electron spectrum. However, this is not the first GRB with these issues. Specifically, this is not the first GRB afterglow for which a hard electron spectrum has been inferred. Different treatments of hard spectra have been put forward in the literature. We have investigated two, and found to give reasonable and physically meaningful results. Finally, the linear polarisation observations [3] are in agreement with a jet break model where the linear polarisation jump is a direct result from the jet break. However, this type of transition in polarisation has not been investigated yet in the literature for energy injection breaks, which might lead to a similar effect.

The results presented here on GRB121024A show that broadband afterglow data from the X-ray to radio allow for a detailed analysis of the characteristic properties of the GRB afterglow synchrotron emission spectrum. Through our extensive data coverage we have been able to constrain the position of all synchrotron breaks, which in turn has allowed us to measure all the micro-physical and dynamical parameters of the GRB afterglow. This information is crucial to study further the GRB afterglow emission processes. Future continual coverage of the GRB afterglow with sensitive telescopes over a wide wavelength range and at multiple epochs will enable us to place strong constraints on the micro-physical parameters for a larger sample of GRBs, and allow us to e.g. investigate the evolution of these parameters.

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References

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