“Hot Edges” in Inverse Opal Structure Enable Efficient CO₂ Electrochemical Reduction and Sensitive in-situ Raman Characterization

Yang Yang a,b, Lukas Ohnoutek c,d, Saira Ajmal a, Xiuzhen Zheng a, Yiqing Feng a, Kejian Li a, Tao Wang a, Yue Deng a, Yangyang Liu a, Xu Dong a, Ventsislav K. Valev c,d*, Liwu Zhang a,b,*,

aDepartment of Environmental Science and Engineering, Fudan University; Shanghai, 200433, P. R. China.

bShanghai Institute of Pollution Control and Ecological Security, Shanghai, 200092, Peoples’ Republic of China

cCentre for Photonics and Photonic Materials, University of Bath, Bath, BA2 7AY, United Kingdom
dCentre for Nanoscience and Nanotechnology, University of Bath, Bath, BA2 7AY, United Kingdom

E-mail: zhanglw@fudan.edu.cn; v.k.valev@bath.ac.uk
Abstract:

Conversion of CO$_2$ into fuels and chemicals via electroreduction has attracted significant interest. Via mesostructure design to tune the electric field distribution in the electrode, it is demonstrated that Cu-In alloy with inverse opal (CI-1-IO) structure provides efficient electrochemical CO$_2$ reduction and allows for sensitive detection of the CO$_2$ reduction intermediates via surface-enhanced Raman scattering. The significant enhancement of Raman signal of the intermediates on CI-1-IO surface can be attributed to electric field enhancement on the “hot edges” of the inverse opal structure. Additionally, a highest CO$_2$ reduction Faradaic efficiency (FE) of 92% (sum of formate and CO) is achieved at -0.6 V vs. RHE on CI-1-IO electrode. The DRIFTS results show that the Cu-In alloy with inverse opal have a faster adsorption kinetic and higher adsorption capacity of CO$_2$. The “hot edges” of the bowl-like structure concentrate electric fields, due to the high curvature and also concentrate K$^+$ on the active sites, which can lower the energy barrier of the CO$_2$ reduction reaction. This research provides new insight into the design of materials for efficient CO$_2$ conversion and the intermediates detection during CO$_2$ reduction processes.
“Hot Edges” in Cu-In monolayer inverse opal structure lead to ultrasensitive Raman signal of the intermediate and efficient electrochemical CO$_2$ reduction performance.
1. Introduction

Electrochemical reduction of carbon dioxide (CO\textsubscript{2}) is a promising way to convert CO\textsubscript{2} to carbon-based fuels and chemicals using renewable electricity.\textsuperscript{1-5} However, its application often suffers from the low conversion efficiency, poor selectivity and the lesser-known intermediates.\textsuperscript{3, 6} Consequently, substantial efforts have been dedicated to designing electrodes with efficient CO\textsubscript{2} reduction activity, such as controlling size and morphology, tuning electrolyte, introducing metal-oxide interfaces and fabricating porous structure etc.\textsuperscript{1, 7-13} By coupling another metal species to form alloy might reconstruct the catalyst interface, which could change the binding energy of intermediates on the surface during CO\textsubscript{2} reduction reaction resulting in controllable selectivity.\textsuperscript{14-21} Yet another method to facilitate CO\textsubscript{2} reduction efficiency and selectivity consists in introducing inverse opal structure.\textsuperscript{22, 23} Surendranath’s group have designed Au and Ag inverse opal thin films with different thickness, which show a significant diminution of HER with increasing porous thin film thickness.\textsuperscript{22, 23} A Cu inverse opal structure derived from Cu\textsubscript{2}O inverse opal was reported to show improved catalytic activity and selectivity of CO\textsubscript{2} electrochemical reduction due to the porous structure.\textsuperscript{24} However, previous work on inverse opal mainly focused on the effects of diffusional gradients within the pores, while many of the specific property of inverse opal structures remain unexplored; such important properties include photonic Bragg resonances, electric field concentrations on the edge of the mesostructure and their potential effects for CO\textsubscript{2} reduction and in-situ measurements.

In-situ measurements have been playing an increasingly important role in the study of the CO\textsubscript{2} reduction mechanism. Such measurements include in-situ Raman spectroscopy, in-situ attenuated total reflectance-fourier transform infrared spectroscopy (ATR-FTIR), vibrational sum-frequency generation spectroscopy (VSFG), and operando extended X-ray absorption fine structure (EXAFS) etc.\textsuperscript{14, 16, 25-30} Among these measurements, in-situ Raman spectroscopy is an efficient and simple method for detecting the intermediates during CO\textsubscript{2} reduction.\textsuperscript{31-39} However,
the Raman signal of the intermediates is very weak, which makes the signal hard to detect.\cite{16,27} Generally, the typical electrochemical Raman cell contains electrodes, an electrolyte layer and a quartz window. The low sensitivity of Raman spectroscopy can be attributed to the limited optical path and mismatch of media.\cite{40} Using an immersion objective with long working distance reduces the loss of Raman signal and improves the detection sensitivity, the spatial resolution, and the electrochemical response.\cite{23,40,41} However, this approach is usually expensive and complex. Nanostructure-based surface enhanced Raman spectroscopy (SERS) is also an effective method for surface analysis,\cite{42} which achieves ultrahigh sensitivity down to a single molecule level by utilizing coinage-metal (such as Au, Ag and Cu) nanostructures.\cite{42,43} Developing electrodes with both good SERS ability and catalytic activity for \textit{in-situ} investigation of the process of CO$_2$ reduction still presents a challenge.

Herein, by depositing a monolayer of inverse opal structured Cu-In alloy, we present a system, rich on “hot edges” for CO$_2$ reduction and \textit{in-situ} Raman investigation. The Raman signal is dramatically enhanced compared to the case when a Cu-In alloy without inverse opal structure is used. Furthermore, the current density, F$\text{E}_{\text{HCOO}}$ and F$\text{E}_{\text{CO}}$ are also significantly improved by introducing the inverse opal structure. The DRIFTS results demonstrate that CI-1-IO with inverse opal structure exhibit a faster adsorption kinetic and higher capacity for CO$_2$ adsorption, which might benefit for the CO$_2$ conversion. The greatly enhanced Raman signal is attributed to local field enhancements, demonstrated by finite-difference time-domain (FDTD) simulations. This Raman signal then enables the \textit{in-situ} study of the CO$_2$ reduction process. The high-curvature structure concentrates electric fields, thus K$^+$ gather on the active sites (“hot edges”). Density functional theory (DFT) calculations further demonstrate that the adsorbed K$^+$ lower the energy barrier of the CO$_2$ reduction reaction (CO$_2$RR), which facilitates the CO$_2$ conversion.
2. Experimental

2.1 Chemicals

Styrene (Aldrich, 99%) and K$_2$S$_2$O$_8$ (AR, Sinopharm Chemical Reagent Co., China) were used to prepare polystyrene (PS) opal templates. Formic acid and Acetic acid standard sample (Sigma), NaHCO$_3$ and NaCO$_3$ (Sigma, 99%) were used for Ion Chromatograph (IC). The other reagents were purchased from Sinopharm Chemical Reagent Co., China. Fluoride-Tin Oxide (FTO) glasses (thickness, 1 mm) were purchased from China Southern Glass Co. Ltd. All chemical reagents are analytical grade and used without any further purification. Deionized (DI) water was used throughout the whole experiment. FTO glasses were sonicated while being immersed in piranha solution, ethanol and DI water in sequence to remove the impurities on the surface.

2.2 Preparation of polystyrene spheres

Monodisperse polystyrene spheres with diameter of about 360 nm were synthesized according to the process reported in the literature. Typically, 100 mL of ultrapure water (18 MΩ/cm$^2$) and 0.6 g styrene (washed by 5% NaOH solution to remove the polymerization inhibitor) were added into a three-necked, round-bottomed flask (250 mL). The mixture was stirred at 500 rpm, while being heated to 70 °C and purged with nitrogen gas with a flow rate of 70 mL/min. Afterwards, the mixture was kept at 70 °C for 20 min, 0.12 g of K$_2$S$_2$O$_8$ was added and kept for 12 h. The oversize and undersize particles of the precursor solutions of polystyrene (PS) microspheres were removed by high-speed centrifugation. A first centrifugation was used to remove the undersize PS microspheres and a second one to remove the oversize PS microspheres. At the end of the process, the average particle diameter of the PS microspheres was approximately 360 nm.

The concentration of PS spheres was diluted to 0.03% (mass fraction) using DI water. Subsequently, the FTO glass was immersed vertically into the PS colloid solution. Then, the solution with the glass was placed in a furnace and kept at 55 °C until the water entirely
evaporated. The PS opal templates self-assembled on the FTO with a close-packing structure, named as FTO-PS.

2.3 Preparation of CuIn film with and without inverse opal structure

Cu-In alloy was prepared by an electrodeposition method reported by Yuan et al. with a minor modification.\textsuperscript{45} Firstly, 5 mmol CuCl\textsubscript{2}.2H\textsubscript{2}O, was added into the solution containing 0.2 mmol triethanolamine and 0.015 mmol sodium citrate. Then, 5 mmol InCl\textsubscript{3}.4H\textsubscript{2}O was added into the solution. Finally, the pH of solution was adjusted to 4 using H\textsubscript{2}SO\textsubscript{4} and NaOH.

2.3.1 Preparation of Cu-In film without inverse opal structure

FTO conductive glass, saturated calomel electrode and platinum foil were used as working electrode, reference electrode and counter electrode, respectively. Electrodeposition was conducted at -1 V (vs. SCE) for 30 min at 30 °C. The sample was named CI-1.

2.3.2 CuIn film with inverse opal structure (CI-1-IO)

For the CI-1-IO preparation, FTO-PS was used as the working electrode and the other electrodeposition conditions were kept the same as for CI-1. After electrodeposition, the sample was dried in air flow, then immersed into tetrahydrofuran (THF) solution for 30 min to remove the PS template. Later, the sample was washed by ethanol and water. We designate this sample: CI-1-IO. Scheme S1 (electronic supplementary information, ESI) illustrates the preparation procedure of polystyrene (PS) opal template and the CI-1-IO sample.

2.4 Characterization

The X-ray diffraction (XRD) patterns were measured using a Bruker D8 Advance X-ray diffractometer with Cu K radiation. The accelerating voltage and applied current were 40 kV and 40 mA, respectively. The general morphology imaging and element mapping of the samples were conducted using a field emission scanning electron microscopy (FESEM, Sigma Zeiss 150, Cal Zeiss Co., Germany) instrument operated at 5 kV and scanning transmission electron
microscopy (STEM, Phenom ProX, at 15 kV) with energy-dispersive spectroscopy (EDS). The photoelectrochemical data was collected using an IVIUM electrochemical workstation. The in-situ diffuse reflectance infrared fourier transform spectroscopy (DRIFTS) spectra were collected on a Shimadzu Tracer-100 FTIR spectrometer equipped with a high-sensitivity mercury cadmium telluride (MCT) detector cooled by liquid nitrogen. A temperature controller was coupled with the DRIFTS chamber (Praying Mantis Kit, Harrick). Samples were placed in a ceramic sample cup (0.35 mm depth, 5 mm i.d.). A series of spectra were recorded every 2 minutes with 100 scans averaged for each spectrum and a resolution of 4 cm⁻¹. The in-situ Raman spectra were measured on an XploRA confocal spectrometer (Jobin Yvon, Horiba Gr, France) with a charge coupled device (CCD) detector. Raman scattering was excited by an external-cavity diode laser (785 nm, with 25% energy) coupled with a 10 × Olympus microscope objective. A holographic notch filter was equipped to filter out the excitation line and a holographic grating (1200 gr/mm) was employed for spectral mapping. A CCD with 1024 × 256 pixels was used to collect the spectra in a resolution of 1.2 cm⁻¹, with 2 accumulations at 2 s acquisition time.

2.5 Electrochemical reduction of CO₂

All electrochemical measurements were carried out on an IVIUM electrochemical workstation in a conventional three-electrode system. A platinum plate and an Ag/AgCl electrode were used as the counter and reference electrodes, respectively. The anode and cathode compartment of electrochemical cell were separated by a Nafion 117 membrane. Before the measurement, CO₂ was bubbled into 50 ml electrolyte of x (x = 0.1, 0.2, 0.5 and 1.0) M KHCO₃ solution for 30 minutes to reach the saturation. The gas products were analyzed by gas chromatography (SHIMADZU, GC-2014) equipped with TDX-1 and HT-POLT Al₂O₃/S columns. A thermal conductivity detector (TCD) was mainly used to detect and quantify H₂, a flame ionization detector (FID) with methanizer was used to analyze and quantify CO by TDX-
1 column, and another FID was used for separating and quantifying CH₄ and other alkane contents with a HT-POLT Al₂O₃/S column. Ar (99.999%) was used as carrier gas. Liquid products were analyzed by Ionic chromatography (IC, Metrohm, 883 Basic IC Plus with Metrosp A SUPP 5-250 column) and ¹H nuclear magnetic resonance (NMR, AVANCE III HD, 500MHz, Bruker) spectrum. In order to confirm the organic products in solution, 0.5 ml liquid production was mixed with 0.1 mL D₂O containing 0.01 uL DMSO as the internal standard for a one-dimensional ¹H NMR spectrum test. Typically, the peak at about 8.3 ppm corresponding to the formic acid (HCOOH) and the peak at ~ 2.6 ppm present the dimethyl sulfoxide (DMSO). To quantify the HCOOH, the solution was diluted 10 times with ultrapure water, then injected into IC. Data are shown in the ESI (Fig. S1-S16).

2.6 Density functional theory (DFT) calculations on Cu₂In surface

DFT calculations were performed with Vienna ab initio Simulation Package (VASP) using the projector augmented wave (PAW) method with the exchange-correlation functional of Perdew-Burke-Ernzeh (PBE) and generalized gradient approximation (GGA).³ ⁴ ⁶ ⁴ ⁸ The (001), (010), (100) and (110) slabs of Cu₂In were constructed by 4 × 4 × 1 supercell and separated by 12 Å of vacuum. The PAW pseudopotentials were used to calculate the interaction between ions and electrons in a plane wave basis set with a cut-off energy of 500 eV and a 4 × 4 × 1 k-points. The Gibbs free energies (ΔG) were calculated at 298.15 K and 1 atm, as outlined below:⁴⁹

\[ G = H - TΔS = E_{DFT} + E_{ZPE} + \int_0^{298.15} C_v \, dT - TΔS \]

Where \( E_{DFT} \) is the optimized total energy, \( E_{ZPE} \) is the zero-point vibrational energy, \( \int_0^{298.15} C_v \, dT \) is the heat capacity, H is enthalpy, T is the temperature and ΔS entropy. The change in ΔG between reaction steps of CO₂ to CO reaction coordinate was calculated from the computational hydrogen electrode model.⁵⁰ Moreover, the binding energy was calculated from DFT-optimized structures as follows: \( E_{Binding} = E_{CO_2} - (E_{Cu_2In} + E_{CO_2}) \), where
$E_{CO_2^*}$ is the energy of system with CO$_2$ proximate to the Cu$_2$In surface, $E_{Cu_2In}$ is the energy of the Cu$_2$In surface (with and without K$^+$ for the respective cases) and $E_{CO_2}$ is the energy of CO$_2$ gas.

2.7 FDTD and COMSOL simulations

For these simulations, the Lumerical FDTD (Finite Difference Time Domain) software was used to calculate the electric fields in the vicinity of the nanostructured surface of the sample CI-1-IO. Simulations were performed for a hexagonal arrangement of the bowl-like structures on the surface. Periodic boundary conditions were applied on the edges of the unit cell. It was assumed that the patterned layer consisted of Cu$_2$In alloy. The permittivity spectra of the alloy were calculated using Bruggeman effective medium approximation. In the model, the structures were irradiated with a linearly polarized broadband optical pulse and the electric field on the surface and in the longitudinal section was monitored.

Simulations of the free electron density on the working electrode with inverse opal structure were performed with the COMSOL Electric Currents module. The simulated system consisted of two electrodes separated with electrolyte. Voltage was applied between the two electrodes. In these simulations, the separation of the electrodes was chosen to be 1 μm and the electrodes were simulated as made of Cu. The electric conductivity of the electrolyte was approximated as 2 S m$^{-1}$. The COMSOL calculations thus provide a good qualitative insight into the behavior of the working electrode during the CO$_2$ reduction process.
3. Results and discussion

3.1. Characterization of catalysts

The detailed steps for synthesizing Cu-In inverse opal (CI-1-IO) is shown in Scheme S1 (described in ESI). The X-ray diffraction (XRD) patterns of FTO, CI-1, and CI-1-IO are shown in Fig. 1a, suggesting the hexagonal-phase Cu$_2$In alloy (JPCDS No. 42-1475, with lattice parameters of $a=b=0.4294$ nm, $c=0.5233$ nm) is obtained. Any peaks from impurities of Cu$_x$O can hardly be seen in the XRD pattern of CI-1-IO. The morphology of the CI-1-IO sample is shown in Fig. 1b, which shows monolayer inverse opal structure (see also Fig. S1). Fig. 1c and d displays the image of an area selected for the energy-dispersive spectroscopy analysis and the total element mapping, respectively. From Fig. 1e to h, it is obvious that the Cu and In elements are distributed on the rim and O, Sn elements from the FTO glass are distributed in the center, which indicates that a monolayer inverse opal structure is formed. The EDX spectra confirm the presence of Cu and In in CI-1 (Fig. S2a) and Cu, In, Sn and O in CI-1-IO (Fig. S2b). The morphology and element mapping of planar CI-1 were also studied and presented in Fig. S1. X-ray photoelectron spectroscopy (XPS) analysis suggested all the three peaks of In$^0$, Cu$^0$ and Cu$_2$O are detected on both CI-1 and CI-1-IO electrodes as shown in Fig. S3. Interestingly, Sn 3d5 spectra can be detected by XPS on the sample CI-1-IO, while it can hardly be observed on the CI-1 suggesting an inverse opal structure is formed on CI-1-IO. From Fig. 1a and S1, it can be seen that for CI-1, the surface of the FTO glass is completely covered by the Cu-In nanoparticles. Therefore, no Sn 3d5 peak is observed for the CI-1 sample. The specific analyses of elements are described below Fig. S3.
**3.2 CO₂ reduction performances of CI-1 and CI-1-IO in H-cell.**

CO₂ reductions on CI-1 and CI-1-IO electrodes were performed at different voltages in 0.1 M KHCO₃ solutions; the results are shown in Fig. 2. It is observed that the FE₉₉₉₉₉₉₉ and FE₉₉₉₉₉₉₀ of CI-1-IO are significantly higher than that of the CI-1 catalyst (Fig. 2a and b), indicating that
inverse opal structure might suppress the H\textsubscript{2} evolution reaction (Fig. 2c) and simultaneously facilitate CO\textsubscript{2} conversion. The highest FE of formate (~ 75\%) is achieved at -0.6 V vs. RHE on CI-1-IO, which is almost 2.7 times higher than that of CI-1 (~ 28\%). Considering that the current density of CI-1-IO is also about 2 times higher than that of CI-I (Fig. 2d), the production rate of formate on CI-1-IO is almost 5.4 times faster than that on CI-1. CI-1-IO catalyst exhibits much higher electrochemical CO\textsubscript{2} reduction FE than CI-1 catalyst, especially at low potentials. At -0.6 V vs. RHE, the highest FE for CO\textsubscript{2} reduction over CI-1-IO reaches ~ 92\%, which is about 2.6 times higher than that of CI-1 (~ 36\%). Therefore, the inverse opal structure not only increases the CO\textsubscript{2} conversion efficiency but also improves the product selectivity. The improved current density and product selectivity cannot be explained by the exposed FTO surface on the CI-1-IO sample, as our experiments using pure FTO showed very low activity for the CO\textsubscript{2} electrochemical reduction (Fig. S7). As the CI-1 and CI-1-IO possess a similar electrode RF of 4.5 and 4.8, the difference in the surface area of the electrodes cannot be the only reason for the improved activity.
Fig. 2. Faradaic efficiency (FE) of a) HCOO\(^-\), b) CO and c) H\(_2\). d) Total current density over Cl-1 and Cl-1-IO catalyst at CO\(_2\)-saturated 0.1 M KHCO\(_3\) solution. The error bar represents the standard deviation of three experiments.

3.3 Mechanism of enhanced CO\(_2\) reduction selectivity and activity

We then further explored the mechanism of enhanced CO\(_2\) reduction selectivity and activity via in situ technique and theoretical calculations. Surendranath and coworkers have designed Au and Ag inverse opal thin films with different thickness, the improved CO\(_2\) reduction activity was attributed to the effects of diffusional gradients within the pores of inverse opal\(^{22-23}\). However, in this work the inverse opal has a monolayer structure (nanobowl array, no three-dimensional continuous pores), the effect of diffusion or mass transportation is thus weak. On one hand, the in-situ DRIFTS results demonstrate that the Cu-In alloy with inverse opal structure exhibits faster adsorption kinetics and higher capacity for CO\(_2\) adsorption,
as shown in Fig. 3. CO₂ adsorption on the sample surface was studied by inducing a CO₂/H₂O mixture to the reaction cell with 10 mL/min flow rate in the dark. The data is collected every two minutes until the CO₂ adsorption peak is stable (about 28 min). The collected in situ DRIFTS spectrum are shown in Fig. 3 (see Fig. S8 for more details). The highest peak at about 2364 cm⁻¹ can be attributed to the CO₂ (gas-phase) adsorption on the sample surface, which show a fast adsorption equilibrium process (about 2 min) for the gas phase. From Fig. S8a, we can see that, after 28 min adsorption, the spectra region of 1000-1800 cm⁻¹ and 2688-3500 cm⁻¹ on CI-1-IO surface are obviously higher than that of CI-1 surface, indicating that the CI-1-IO with inverse opal structure have a higher adsorption ability than that of CI-1. The broad band at 2688-3500 cm⁻¹ correspond to the increased hydrogen bonding between the CO₂ adsorption products and hydroxyl groups on the Cu-In alloy surface. The spectra region from 1000 to 1800 cm⁻¹ can be assigned to the adsorbed carbonate and bicarbonate, which are shown in detail in Fig. 3c and 3d (also shown in Fig. S8b and c), respectively. Specifically, the band at about 1227 and 1686 may be attributed to the δ₄ (COH) and ν₂ (OCO)₄ of vibrational modes of the adsorbed bicarbonate. It has been reported that the carbonate band usually appear at 1200-1700 cm⁻¹ and often observed as broad overlapping bands. The band at about 1566 cm⁻¹ may correspond to the bidentate form of absorbed carbonate, and the peak at 1057 and 1383 cm⁻¹ is associated with the monodentate form of carbonate, usually considered as ν₁ stretch. Notably, the intensity of the peaks of carbonate and bicarbonate increase with the increasing of CO₂ flow time shown in Fig. 3c and d, and after 28 min the intensity of the peaks of CI-1-IO is much stronger than that of CI-1. Specifically, the intensity changes of 1383 and 1227 cm⁻¹ are summarized in Fig. S3b. Additionally, the CO₂ adsorption performance on pre-wetted electrode surface in closed cell is also studied as well (shown in Fig. S9). Fig. S9 shows that the adsorbed gas-phase CO₂ (about 2366 cm⁻¹) was saturated for both CI-1 and CI-1-IO after 3 mins. Meanwhile, the CI-1-IO sample shows a higher adsorption capacity of gas-phase CO₂ than CI-1. More details are discussed in ESI. The strong adsorption intensity and fast adsorption
equilibrium process demonstrate CI-1-IO have a faster adsorption kinetic and higher capacity for CO\textsubscript{2} adsorption, which would benefit for the CO\textsubscript{2} conversion.

\textbf{Fig. 3.} \textit{In-situ} DRIFTS spectra of CO\textsubscript{2}/H\textsubscript{2}O adsorbed on CI-1 and CI-1-IO surface. a) Whole spectra from 750-4500 cm\textsuperscript{-1}, 0 to 28 min. b) Intensity of specific peaks (1383 and 1227 cm\textsuperscript{-1}) of CI-1 and CI-1-IO. Spectra regions (1000-1800 cm\textsuperscript{-1}) of adsorbed carbonate on c) CI-1 and d) CI-1-IO surface.

On the other hand, the high-curvature structures in inverse opal can concentrate electric field, which affects the ion concentrations near the surface. To understand the mechanism of improved CO\textsubscript{2} reduction on CI-1-IO, we detected the local relative K\textsuperscript{+} concentration adsorbed on the electrodes by ion chromatography (IC) spectrometry. The concentration of K\textsuperscript{+} adsorbed on the electrode surface was measured by the reported method and shown in Fig. S10.\textsuperscript{3} It was found that the K\textsuperscript{+} concentration when CI-1-IO was used (50.3 mg L\textsuperscript{-1}) was about 2 times higher than in the case of CI-1 (24.5 mg L\textsuperscript{-1}).
The influence of K\(^+\) on the CO\(_2\)RR was further studied via Ab initio molecular dynamics simulations. The Gibbs free energy of diagrams obtained from density functional theory (DFT) calculations on Cu\(_2\)In surface models of different facets at 298.15K, 1 atom and 0 V vs. RHE are shown in Fig. 4. The results indicate that the adsorbed K\(^+\) ions can lower the thermodynamic energy barrier for reaction for all facets. On Cu\(_2\)In (001) facet (Fig. 4a), the adsorbed K\(^+\) stabilizes the COOH\(^*\) and HCOOH\(^*\) intermediates by 1.21 and 1.82 eV, respectively. For Cu\(_2\)In (100) and (110) facet, the energy barrier of rate-determining COOH\(^*\) is lowered by 0.75 and 0.73 eV with the presence of K\(^+\), respectively (Fig. 4b and c). Meanwhile, Fig. 4b and c show that adsorbed K\(^+\) stabilizes the HCOOH\(^*\) by 0.48 and 0.6 eV for the Cu\(_2\)In (100) and (110) surface, respectively. On the Cu\(_2\)In (010) surface (Fig. 4d), the adsorbed K\(^+\) lowers the energy of COOH\(^*\) and HCOOH\(^*\) from 1.09 to 0.48 eV and 0.8 to 0.25, respectively. Simultaneously, the adsorbed K\(^+\) ions also show an ability to lower the thermodynamic energy barrier for the production of CO for all facets (seeing Fig. S11) and the optimized structures for different Cu\(_2\)In facets with and without K\(^+\) are shown in Fig. S12. Overall, the results show that the adsorbed K\(^+\) might lower the reaction barrier of COOH\(^*\), HCOOH\(^*\) and CO\(^*\) intermediates which demonstrates that the adsorbed cations are beneficial for the CO\(_2\) reduction.

Based on the results discussed above, it can be suggested that K\(^+\) ions locally concentrated at the active sites enhance the CO\(_2\)RR. The high-curvature structures in an inverse opal can concentrate electric field, which affects the ion concentrations near the surface. Here, a finite-element numerical calculation is employed to study the enhanced field intensification caused by the inverse opal structure. The computed electric field near the sharp edges of the bowl structure is shown in Fig. 4e. The free electron density distribution on the edges of bowl structure is shown as a color map. The electrostatic field distribution around the edge of the bowl-like structure on the electrode is shown as a group of arrows, where the size and direction of each arrow represent the magnitude and the direction of the electric field. It can be seen that the edges of the bowl structure concentrate electrons. Additionally, a locally enhanced
The electrostatic field is generated by the locally concentrated free electron density on the surface of the electrode and points towards the locally concentrated free electron density (arrows in the inset in Fig. 4e). The results can be explained by migration of free electrons to the regions of the highest curvature on a charged metallic electrode which is a consequence of electrostatic repulsion.

![Diagram](image)

**Fig. 4.** Thermodynamic barriers for $\mathrm{CO}_2$ to HCOOH reduction reaction on Cu$_2$In surface with and without K$^+$. Gibbs free energy of electroreduction of $\mathrm{CO}_2$ to HCOOH on Cu$_2$In a) (001), b)
(100), c (110) and d) (010) facets in the presence of adsorbed K⁺ and without K⁺. e) Computed electron density and the electric field near the surface of the CI-1-IO electrode. The color map represents the normalized free electron density distributed on the surface of the material. The direction and the magnitude of the electrostatic field around the electrode is represented by the white arrows in the inset.

3.4 **In-situ** Raman investigation

Raman signal can also benefit from the inverse opal structure, which enhances local electric fields. It is thus of interest to use this structure for *in-situ* investigation of the reduction process of CO₂ via surfaced enhanced Raman spectroscopy. Fig. 5a-c show the *in-situ* Raman spectra obtained in a tailored cell (shown in Fig. S13) utilizing an electrolyte containing 7 ml CO₂ saturated 0.1 M KHCO₃ over the CI-1 and CI-1-IO electrodes at -0.7 V vs. RHE. The chemical composition of the prepared electrodes and their transformations during CO₂ electroreduction were elucidated by Raman spectroscopy (Fig. 5). In the microscopic image of CI-1-IO (Fig. 5a), we can see some bright spots on the electrode surface caused by the inverse opal structure. These are consistent with surface-enhancement of the Raman scattering signal from the chemical compounds on the surface. In Fig. 5b, it can be seen that the Raman signal of the reaction intermediates can hardly be identified on the CI-1 electrode without inverse opal structure, while the signal is much clearer when the CI-1-IO electrode is used, as shown in Fig. 5c. The Cu-In alloy with inverse opal structure has a SERS effect, which makes it easier to measure the Raman spectra. Before applying external voltage (at time = 0 s), the peaks at 219, 526 and 617 cm⁻¹ can be assigned to Cu₂O.²³ Within 60 s of applying potential difference, the peaks (at 526 and 617 cm⁻¹) attributed to Cu₂O disappear, which indicates that the catalyst is in the metallic state. Meanwhile, the CO-related peak at 2075 cm⁻¹ (CO stretching) appears. A strong signal at 1065 cm⁻¹ is also observed, which can be assigned to the stretching mode of the adsorbed carbonate in the interfacial region.³¹, ³⁵, ⁵⁸ This signal was not observed in recent
reports.\textsuperscript{16, 23, 26, 33, 59} Up to now, it is poorly understood that how the carbonate form and whether it affect CO\textsubscript{2} reduction. One view is that carbonate adsorbed on the metal surface is thorough ionic bonds. Such a bond is formed by ligand exchange of HCO\textsubscript{3}\textsuperscript{-} with adsorbed water/hydroxyl followed by deprotonation, which does not participate in the CO\textsubscript{2} reduction reaction.\textsuperscript{6, 35} Another viewpoint thought the carbonate is formed by activation of CO\textsubscript{2} on the surface oxygen/hydroxyl sites.\textsuperscript{6} There is another proposition (proved by the ATR-FTIR spectra) that CO\textsubscript{2} couple with CO\textsubscript{2}\textsuperscript{-} intermediate firstly, then reduced by another electron resulting in CO and CO\textsubscript{3}\textsuperscript{2-} formation.\textsuperscript{27} The broad band at about 2900 cm\textsuperscript{-1} can be assigned to the C-H stretching of the formate species.\textsuperscript{58, 60}

During the study time from 60 to 600 s, the signals of carbonate (1065 cm\textsuperscript{-1}) and CO (2075 cm\textsuperscript{-1}) stretching are steady, indicating that a balance of the reduction reaction is reached. When the voltage is removed for 3 min (the first red curve in Fig. 5c), the intensity of the carbonate and CO stretching signal decreases. The Cu\textsubscript{2}O peaks appear again, while the carbonate and CO stretching signal decrease significantly after removing the voltage for 10 min. After removing the voltage for 30 min, the Raman spectrum returns to the original shape before reaction. Neither the signal from carbonate nor from CO stretching is observed. However, when the voltage is applied again for 30 s, the signal of Cu\textsubscript{2}O disappears again, while the carbonate and the CO stretching mode peaks can be observed. After removing the potential for 30 min, the Raman spectrum is similar to the first round, indicating the repeatability of the Raman study. These operando Raman spectra demonstrate that the Cu\textsubscript{2}O is rapidly reduced to the metallic state upon the application of voltage and then the CO\textsubscript{2} reduction reaction occurs on the metal surface. Interestingly, the reduced Cu reoxidizes to Cu\textsubscript{2}O after removing the applied voltage. Both the CO\textsubscript{3}\textsuperscript{2-} and CO intermediates (CO stretching) can be observed during the CO\textsubscript{2} reduction in the Raman spectra. The inverse opal structure enables more sensitive \textit{in-situ} Raman study of the reaction process during CO\textsubscript{2}RR.
**Fig. 5.** a) Microscopic image of CI-1-IO. *In-situ* Raman spectra of b) CI-1 and c) CI-1-IO during CO₂ reduction at -0.7 V vs. RHE in 0.1 M KHCO₃.

To further understand the influence of the inverse opal structure on the electric field at the surface of the electrode, finite-difference time-domain (FDTD) method was used to calculate the electric field distribution on the Cu-In inverse opal structure when irradiated with light. Fig. 6a and b show the distribution of electric field (E-field) on the surface and the longitudinal profile of the field when the Cu-In monolayer with inverse opal structure is irradiated with 785 nm light. It can clearly be seen that the center of the cavity, the top and the sharp edges of the bowl structure show an enhanced E-field. Specifically, on the sharp edges of the bowl, the magnitude of the E-field is 2-4 times higher. It has been reported that a small increase in the local field can produce large enhancements in Raman scattering (the enhancement scales roughly as E⁴). The enhanced Raman signal can also benefit from the improved production rate of CO and formate on CI-1-IO, which is 4.8 and 3.4 times faster than on CI-1 at -0.7 V vs. RHE. However, the Raman signal is improved significantly more than this, as seen in Fig. 5b. The Raman signal of the intermediates can hardly be observed on CI-1. Hence, the enhancement of Raman scattering can be mainly attributed to the enhanced electric field on the surface of the...
inverse opal structure. The “hot edges” presents as an active site is also supported by previous report by Morimoto et al., who use open-loop electric potential microscopy (OL-EPM) prove that the edge site of porous Zn is the reactive site for CO$_2$ reduction.$^{62}$ Additionally, their OL-EPM results show that the higher current concentration at the edge site is attribute to the increased electric field strength.

![Diagram](image)

**Fig. 6.** Electric field distribution simulated for Cu-In inverse opal structure, when irradiated with 785 nm light, by a finite-difference time-domain (FDTD) method. The direction of light polarization is along x. In a), mapping of the electric field distribution on the surface. In b), mapping of the longitudinal profile of the electric field distribution.
4. Conclusions

In summary, we have successfully prepared a Cu-In alloy electrode with inverse opal structure possessing plenty of “hot edges” for CO$_2$ reduction and showed that it can be used *in-situ* to improve the sensitivity of Raman spectroscopy. The structure not only exhibits higher efficiency of CO$_2$ reduction, but also dramatically enhances the Raman spectra compared to Cu-In alloy without inverse opal structure. DRIFTS results display that the inverse opal structure has a faster adsorption kinetic and higher capacity for CO$_2$ adsorption, which would benefit for the CO$_2$ conversion. The Cu-In alloy electrode with inverse opal structure also concentrates electric fields due to the high curvature structure (COMSOL simulation), and thus concentrates K$^+$ on the active sites, which lowers the thermodynamic energy barrier of the CO$_2$ reaction (DFT calculation) and reduces the mean square displacement of CO$_2$ relative to the surface. The characteristic peaks of the CO intermediate, carbonate and formate were observed on the CI-1-IO electrode. However, these peaks can hardly be observed on the CI-1 electrode. The enhanced Raman signal can be attributed to the inverse opal structure, which enhances the electric field on the sharp edges of the electrode (FDTD simulation). This work introduces “hot edges” structure as promising strategy for improving efficiency and selectivity of CO$_2$RR, this structure also enables sensitive *in-situ* detection of intermediates via SERS and thus provides a way to elucidate the exact mechanism of the electrochemical CO$_2$RR.
Conflicts of interest

There are no conflicts to declare.

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41. Y. L. Deng and B. S. Yeo, ACS Cata., 2017, 7, 7873-7889.
43. S. Nie and S. R. Emory, Science, 1997, 275, 1102-1106.