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# 1 HEAT TRANSFER IN TURBOCHARGER TURBINES UNDER STEADY, PULSATING 2 AND TRANSIENT CONDITIONS

3

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## 8 ABSTRACT

9 Heat transfer is significant in turbochargers and a number of mathematical models have been  
10 proposed to account for the heat transfer, however these have predominantly been validated under  
11 steady flow conditions. A variable geometry turbocharger from a 2.2L Diesel engine was studied,  
12 both on gas stand and on-engine, under steady and transient conditions. The results showed that  
13 heat transfer accounts for at least 20% of total enthalpy change in the turbine and significantly more  
14 at lower mechanical powers. A convective heat transfer correlation was derived from experimental  
15 measurements to account for heat transfer between the gases and the turbine housing and proved  
16 consistent with those published from other researchers. This relationship was subsequently shown  
17 to be consistent between engine and gas stand operation: using this correlation in a 1D gas dynamics  
18 simulation reduced the turbine outlet temperature error from 33°C to 3°C. Using the model under  
19 transient conditions highlighted the effect of housing thermal inertia. The peak transient heat flow  
20 was strongly linked to the dynamics of the turbine inlet temperature: for all increases, the peak heat  
21 flow was higher than under thermally stable conditions due to colder housing. For all decreases in  
22 gas temperature, the peak heat flow was lower and for temperature drops of more than 100°C the  
23 heat flow was reversed during the transient.

24 **Keywords:** Turbocharger, Heat transfer, Transient, Thermal modelling

## 25 **1 INTRODUCTION**

26 Turbocharging internal combustion engines is set to increase rapidly as this is a key technology to  
27 deliver fuel economy savings for both Diesel and spark ignition engines [1]. Using a compressor to  
28 provide higher air flows to an internal combustion engine increases the power density and allows  
29 smaller engines to be used in more high power applications, reducing overall weight and friction.  
30 The matching of a turbocharger with an internal combustion engine is a crucial step in the  
31 development process and relies on simulation of the engine air path system. In these models,  
32 turbochargers are represented by characteristic maps, which are defined from measurements of  
33 pressure ratio, shaft speed, mass flow and isentropic efficiency taken from a gas stand. Whilst the  
34 mass flow, pressure ratio, and speed can be measured directly, the efficiency has to be calculated  
35 from measured gas temperatures. For both turbine and compressor, enthalpy changes in the  
36 working fluids are equated to work changes during the characterisation process<sup>1</sup>. Any heat transfer  
37 affecting these gas temperature measurements will cause errors in the characterisation process.  
38 Conversely, when the characteristic maps are subsequently used in engine simulations to predict  
39 engine performance; if heat transfers are ignored then a poor prediction of gas temperatures for  
40 inter-cooling and after-treatment will arise. Consequently there is a two-fold interest in  
41 understanding and modelling heat transfer in turbochargers:

- 42 1. To improve the accuracy of work transfer measurements during characterisation.
- 43 2. To improve the prediction of gas temperatures in engine simulations.

44 Current practice ignores heat transfers and limits investigations to operating conditions where heat  
45 transfer are small compared to work transfers; these conditions prevail for the compressor at higher  
46 turbocharger speeds but heat transfer is always significant in the turbine. Parametric curve fitting  
47 techniques are then used to extrapolate to the lower speed region [2].

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<sup>1</sup> Some specialist facilities use a turbine dynamometer to measure turbine work directly, however these rarely used for automotive turbochargers in industrial applications.

48 This work focuses on heat transfer in the turbine which represents the principal heat source for  
49 turbocharger heat transfer and strongly affects the gas temperature entering after-treatment  
50 systems. In particular, this paper aims to assess the applicability of gas-stand derived heat transfer  
51 models to on-engine conditions where flows are hotter, pulsating and highly transient.

## 52 **2 BACKGROUND**

53 A number of studies into heat transfer in turbochargers have been presented over the past 15 years.  
54 The first studies focussed on quantifying the effects of heat transfer on steady flow gas stands by  
55 comparing the work transfers that would be measured based on temperature changes for different  
56 turbine inlet temperatures [3-8]. Cormerais et al. [4] presented the most extreme changes in  
57 operating conditions, varying turbine inlet temperature from 50°C to 500°C with a thermally  
58 insulated turbocharger and observed up to 15%points change in apparent compressor efficiency .  
59 Baines et al. [7] measured losses of 700W at 250°C turbine inlet gas temperature (TIT) which is  
60 considerably lower than the 2.7kW measured for a similar turbocharger by Aghaali and Angstrom  
61 with turbine inlet temperatures ranging 620-850°C [8]. Baines et al. [7] also estimated heat transfer  
62 to ambient as 25% of total turbine heat transfer, however at 700°C TIT, where temperature  
63 gradients to ambient were much higher, Shaaban [5] estimated this at 70%.

64 A number of modelling approaches have been used ranging from 3D conjugate heat transfer, giving  
65 a detailed insight to the heat transfer processes [9, 10], to simple 1D models for use with engine  
66 simulations. The most basic approach adopted to improve the correlation of engine models to  
67 experimental data consists of empirically adapting or *correcting* turbine maps using efficiency  
68 multipliers [8, 11]. This approach is typically parameterized to estimate heat energy directly using an  
69 exponential function that decays with increasing mass flow or turbine power and is tuned to match  
70 measured data from an engine or vehicle dynamometer. Whilst this approach can improve the  
71 accuracy of engine models, it is not predictive and alternative models have been proposed.

72 In practice heat transfer will occur through the turbocharger stage [12], however a common  
 73 assumption in 1D models assumes that heat transfer and work transfer occur independently [13-15];  
 74 this is represented schematically on enthalpy-entropy diagrams in figure 1. The actual processes  
 75 undergone by the gases are shown between states 1-2 and 3-4 for compressor and turbine  
 76 respectively. The split of work and heat transfer is shown by the intermediate states 1', 2', 3' and 4'  
 77 such that flow through the turbine is composed of the following stages:

- 78 1. A heating or cooling at constant pressure (processes 1-1' and 3'-3),
- 79 2. An adiabatic compression/expansion (processes 1'-2' and 3'-4')
- 80 3. A heating or cooling at constant pressure (processes 2'-2 and 4'-4)

81 Based on this analysis it is obvious that any measurement of temperature change across the turbine  
 82 or compressor will include both the work and heat transfers, and that any estimate of work based on  
 83 the total enthalpy change will include an error equal to the net heat transfer (equation 1).

$$\Delta h_{act} = \Delta h_{work} + q_b + q_a \quad 1$$

84

85 The isentropic efficiencies used in engine simulation codes are described for compressor and turbine  
 86 in equations 2 and 3 respectively. These equations describe the transitions between 1'-2' and 3'-4'.

$$\eta_{s,c} = \frac{\Delta h_{s',c}}{\Delta h_{work,c}} = \frac{c_{p,c} \left[ T_{01'} \left( \left( \frac{P_{02}}{P_{01}} \right)^{\frac{\gamma-1}{\gamma}} - 1 \right) \right]}{\Delta h_{act,c} - q_{b,c} - q_{a,c}} \quad 2$$

$$\eta_{s,t} = \frac{\Delta h_{work,t}}{\Delta h_{s',t}} = \frac{\Delta h_{act,t} - q_{b,t} - q_{a,t}}{c_{p,t} \left[ T_{03'} \left( 1 - \left( \frac{P_4}{P_{03}} \right)^{\frac{\gamma-1}{\gamma}} \right) \right]} \quad 3$$

87

88 In equations 2 and 3 it is common to define efficiencies using total conditions at points 1, 2 and 3  
89 (and hence 1', 2' and 3') and static conditions at point 4 (and 4'). For clarity, these distinctions have  
90 been omitted from figure 1.

91 The major issue that arises in applying equations 2 and 3 is that it is not possible to directly measure  
92  $T_1$ ,  $T_2$ ,  $T_3$ ' and  $T_4$ ' because they are not well defined spatially within the turbocharger. Consequently,  
93 for industrial mapping, operation is assumed to be adiabatic, i.e.  $q_a=q_b=0$ ,  $T_1=T_1'$ ;  $T_2=T_2'$ ,  $T_3=T_3'$  and  
94  $T_4=T_4'$ . This assumption holds for a compressor operating at higher shaft speeds where the heat  
95 transfer is small compared to the work transfer [16]. On the turbine side, the condition of adiabatic  
96 operation can only be achieved in special laboratory conditions and commonly turbine work is  
97 estimated either through compressor enthalpy rise or using a turbine dynamometer [17].

98 The 3D conjugate heat transfer modelling undertaken by Bohn et al [9] showed that heat transfers  
99 between the working fluids and the housing could occur in either direction and could change  
100 direction as the flow passed through the rotor and diffuser depending on the magnitude of  
101 temperature change due to compression or expansion. To capture this in a simplified model, the full  
102 problem described by figure 1 should be considered where heat transfers can occur both before and  
103 after the compression and expansion processes. However, most authors [14, 18, 19] prefer to group  
104 all heat transfers after the compression in the compressor or before the expansion in the turbine:  
105 i.e. in figure 1 (a)  $q_b = 0$  and in figure 1 (b)  $q_a = 0$ . This simpler approach stems from a limitation in the  
106 parameterisation method. This is performed either by comparing hot operation of the turbocharger  
107 with special conditions where temperature gradients are minimised by matching  $T_2$  and  $T_3$ <sup>2</sup>, or by  
108 using the turbocharger bearing housing as a heat flux probe [15]. In both cases further assumptions  
109 are required for separating the heat flows before and after work transfers [20] and these are  
110 deemed not to provide any further accuracy benefits over lumping all heat transfers into a single  
111 process. The convective heat transfer between the working fluid and the housing within the turbine

---

<sup>2</sup> The turbocharger cooling fluids (oil and water if present) are also matched to the compressor outlet and turbine inlet gas temperatures.

112 and compressor housing is always modelled by assuming or adapting convective correlations for  
 113 flows in pipes such as Dittus-Boelter or Seider-Tate [21]. A number of correlations proposed in the  
 114 literature are presented in table 1. It is difficult to compare these correlations in equation forms  
 115 because of differences in defining the characteristic lengths. Therefore a graphical representation is  
 116 given in the results section of this paper (figure 12).

117 **Table 1: Comparison of internal convective heat transfer correlations for turbines**

Authors	Source	Correlation	Characteristic Length	Constants		
				a	b	c
Baines et al. [7]	Gas stand	$Nu = aRe^bPr^c$	$L_{volute}$	0.032	0.7	0.43
Cormerais [18]			$D_{inlet}$	0.14	0.75	1/3
Reyes-Belmonte [22]	Gas Stand	$Nu = aRe^bPr^{1/3} \left(\frac{\mu_{bulk}}{\mu_{skin}}\right)^{0.14} F$ <i>where</i> $F = 1 + 0.9756 \left(\frac{\frac{D_{inlet}}{\eta_{max}}}{\frac{(L_{volute})^2}{4D_{inlet}}}\right)^{0.76}$	$\frac{(L_{volute})^2}{4D_{inlet}}$	1.07	0.57	1/3
				5.34	0.48	1/3
				0.101	0.84	1/3
Romagnoli and Martinez-Botas [19]	Theory	$Nu = aRe^bPr^c$	$\frac{D_{inlet}}{2}$	0.046	0.8	0.4

118

119 The heat transfer models have been shown to improve the accuracy of turbine outlet temperature  
 120 prediction from an over prediction of 20-40°C to within ±10°C [16]. However, no direct comparison  
 121 has been made for the same device between gas stand and engine operation. In this paper, an  
 122 investigation with the same turbocharger and crucially the same instrumentation was conducted in  
 123 both environments.

## 124 3 MODELLING AND DATA ANALYSIS

### 125 3.1 Total Heat Transfer

126 An overview of the heat and work flows inside the turbocharger is shown in figure 2. By applying the  
127 conservation of energy, the change in enthalpy in the turbine can be related to the work and heat  
128 transfer rates using equation 4, with  $T_{0i}$  the stagnation or total temperatures.

$$\dot{W}_t + \dot{Q}_{b,t} + \dot{Q}_{a,t} = \dot{m}_t c_{p,t} (T_{03} - T_{04}) \quad 4$$

129

130 Where the turbine work transfer rate can be derived from a power balance on the shaft (equation  
131 5).

$$\dot{W}_t = \dot{W}_c + \dot{W}_f \quad 5$$

132

133 The compressor work transfer rate is estimated using equation 6; this effectively ignores heat  
134 transfers on the compressor side. This will cause errors, notably at low speeds and a full analysis of  
135 the uncertainties caused by this assumption are given in section 4.4. The friction work was estimated  
136 using the model developed by Serrano et al [23] , summarized by equation 7.

$$\dot{W}_c = \dot{m}_c c_{p,c} (T_{02} - T_{01}) \quad 6$$

$$\dot{W}_f = C_{fr} N_t^2 \quad 7$$

137 Combining equations 4 with equations 5-7 and rearranging yields the expression for total heat  
138 transfer from the gas to turbine housing:

$$\begin{aligned} \dot{Q}_{G/T} &= \dot{Q}_{b,t} + \dot{Q}_{a,t} \\ &= \dot{m}_t c_{p,t} (T_{03} - T_{04}) - C_{fr} N_t^2 \\ &\quad - \dot{m}_c c_{p,c} (T_{02} - T_{01}) \end{aligned} \quad 8$$

139



### 140 3.2 Heat Transfer model

141 A simplified heat transfer model was used based on similar approaches found in the literature  
142 [4,5,7,13,14,19] (figure 2). The model combines two thermal nodes (compressor and turbine  
143 housing), linked via conduction through the bearing housing. Heat transfer between the gases and  
144 housings can occur both before and after the compression/expansion processes which is important  
145 because of the different temperature gradients between gas and wall.

146 The focus of this paper remains on the heat transfer between the exhaust gases and the turbine  
147 node. Undertaking an energy balance on this node yields equation 9; the heat transfer model aims  
148 to determine each of the terms on the right hand side.

$$m_T c_{p,T} \frac{dT_T}{dt} = \dot{Q}_{b,T} + \dot{Q}_{a,T} - \dot{Q}_{T/B} - \dot{Q}_{T,rad} - \dot{Q}_{T,conv} \quad 9$$

149

150 To eliminate  $\dot{Q}_{T,rad}$  and  $\dot{Q}_{T,conv}$  from equation 9, a measured turbine housing temperature was  
151 used which avoids the uncertainties in modelling external heat transfer, most notably with respect  
152 to external air flows which strongly affect the convection term [24], are highly specific to different  
153 installations (gas stand, engine dynamometer, in-vehicle) and difficult to capture without a full 3D  
154 simulation.

155 The flow path inside the turbine is highly complex with variations in section, flow rates and  
156 convective area. For the turbine, the tongue could be approximated to a short pipe of constant  
157 diameter, however the scroll has a gradually reducing diameter and mass flow rate as gas enters the  
158 stator and rotor flow passages. The flow is then combined in the diffuser, which may once again be  
159 approximated as constant diameter pipe. From a heat transfer perspective, this means that the large  
160 spatial variations in flow conditions will result in a wide range of local Reynolds numbers that would  
161 be difficult to validate experimentally.

162 In this simplified model, the turbine is considered as two pipes of constant diameter, with an  
163 adiabatic expansion between them. The heat transfer in the pipes is calculated using Newton's law  
164 of cooling (equations 10 and 11). The total wetted area, ( $A_T = A_{b,T} + A_{a,T}$ ) can be determined  
165 from part geometry. The breakdown of area pre- and post- compression in this paper is assumed to  
166 be 85% of total area before expansion and 15% after, which has been determined based on a  
167 qualitative assessment of static temperature drop through the turbine. Whilst a more rigorous  
168 approach to determining this breakdown in area could be desirable, previous work on heat flows in  
169 compressor housings showed that this breakdown in heat flow only becomes significant if there are  
170 large pressure changes in the device [20]. Therefore the arbitrary assignment of distribution in this  
171 work is deemed sufficient. Heat flows presented in the subsequent sections of this work consider the  
172 total heat transfer over the complete turbine. In this way a spatially averaged Reynolds number is  
173 defined for the whole turbine stage, acting over the total heat transfer area. To account for the  
174 geometry of the device, the constants  $a_1$  and  $a_2$  of the Seider-Tate convection correlation (equation  
175 12) [21, 25] were determined empirically based on measured gas and wall temperatures.

$$\dot{Q}_{b,T} = h_{b,T} A_{b,T} (T_3 - T_T) \quad 10$$

$$\dot{Q}_{a,T} = h_{a,T} A_{a,T} (T_{4'} - T_T) \quad 11$$

with  $A_T = A_{a,T} + A_{b,T}$

$$Nu_t = \frac{h_{x,T} D_{T,inlet}}{k_G} = c_1 Re_t^{c_2} Pr^{1/3} \left( \frac{\mu_{bulk}}{\mu_{skin}} \right)^{0.14} \quad 12$$

176

177 It is vitally important the definition of characteristic length  $D_T$  and effective heat transfer area  $A_{x,T}$  be  
178 provided with the convection correlation parameters  $c_1$  and  $c_2$  in equation 12 as many dimensions  
179 could be considered for this. Here, the inlet and outlet diameters are defined as characteristic  
180 lengths whilst the internal heat transfer area is the total area as calculated from part geometry.

### 181 3.3 Work transfer and Gas Dynamic Model

182 The model assumes that work and heat transfer occur independently. The enthalpy change due to  
183 heat transfer is captured by the model described above whilst the work transfer will be captured by  
184 an isentropic efficiency term that represents the isentropic efficiency that would be observed  
185 experimentally if no heat transfer were present. These were derived from the isentropic efficiencies  
186 derived from the gas stand measurements,  $\eta_{gas\ stand}$  (equation 13).

$$\eta_{gas\ stand} = \frac{T_{02} - T_{01}}{T_{03} - T_{4s}} \quad 13$$

187

188 However this efficiency term will not account for mechanical losses in the turbocharger bearing as it  
189 is based on the apparent work transfer in the compressor. Mechanical losses were estimated using  
190 the friction model developed by Serrano et al. [23] (see equation 7). Equation 14 then uses this  
191 estimate and the compressor work (equation 6) to calculate mechanical efficiency.

$$\eta_{mech} = \frac{W_c}{W_f + W_c} \quad 14$$

192

193 Equation 15 can then be used to calculate an actual turbine efficiency.

$$\eta_s = \frac{\eta_{gas\ stand}}{\eta_{mech}} \quad 15$$

194

195 To account for the behaviour of the turbine under pulsating flow conditions such as those observed  
196 on-engine, a mean line turbine model was used [26]. This represents the turbine as a series of two  
197 orifices and an internal volume and was used to calculate instantaneous Reynolds numbers from  
198 pulsating pressure measurements.

## 199 4 EXPERIMENTAL APPROACH

### 200 4.1 Turbocharger description

201 The turbocharger used in this study was from a 2.2L automotive Diesel engine, with turbine and  
202 compressor wheel diameters of 43mm and 49mm respectively. The turbine side included variable  
203 guide vanes and cooling was provided by engine lubricating oil.

204 K-type thermocouples were installed to measure fluid and metal temperatures. At each gas inlet and  
205 outlet port, three thermocouples were installed with 0.5, 0.3 and 0.15 diameter protrusions into the  
206 flow (figure 3a). These depths were chosen arbitrarily to give a radial temperature distribution. The  
207 number of sensors that could be installed was limited by space constraints within the engine  
208 components while a greater number of sensors would improve the knowledge of temperature  
209 distribution. The sensors were installed through the housing of the turbocharger and therefore as  
210 close as possible to the compression and expansion processes thus minimising heat losses between  
211 the measurements. In this way, a distribution of temperature is captured at the four gas ports of the  
212 turbocharger and is able to capture to a degree the non-homogeneous temperature that exists at  
213 these ports [26]. Crucially the instrumentation between gas stand and on-engine remains constant  
214 allowing for direct comparison between the two configurations. Additional thermocouples were  
215 installed in the bearing, turbine and compressor housings in order to estimate a bulk metal  
216 temperature (figure 3b and c).

### 217 4.2 Gas stand test facility

218 A schematic of the gas stand facility is given in figure 4: the turbine is supplied with hot compressed  
219 air from a screw compressor and electrical heating system. The flow through the turbine is  
220 controlled through an electric valve and measured using a thermal flow meter before being  
221 thermally conditioned by two electric heaters. The turbine drives the compressor and flow through  
222 the compressor is controlled by a second electric valve, and measured using a second thermal flow  
223 meter. Temperatures and pressures are measured at the inlet and outlet of both devices using k-

224 type thermocouples and piezo-resistive sensors respectively. Additional k-type temperature sensors  
225 integral to the gas stand facility were also available for these experiments. The turbocharger is  
226 lubricated using a dedicated oil supply and conditioning system, ensuring oil temperature remains  
227 above 70°C during all experiments.

228 Turbine maps were measured under thermally stable conditions for two turbine inlet temperatures  
229 (100°C and 500°C) and three variable geometry turbine (VGT) positions (20%, 50% and 80%). The  
230 operating points are shown on the compressor and turbine maps in figure 5. Corrected flow  
231 conditions for compressor and turbine were obtained from equations 16 and 17 respectively.

$$\dot{m}_{c,corr} = \dot{m}_c \frac{\sqrt{\frac{T_{01}}{298}}}{\frac{P_{01}}{1}} \quad 16$$

$$\dot{m}_{t,corr} = \dot{m}_t \frac{\sqrt{\frac{T_{02}}{288}}}{\frac{P_{03}}{1.01325}} \quad 17$$

232

### 233 4.3 Engine test facility

234 A 2.2L Diesel engine was installed on a transient AC dynamometer and was used to control  
235 turbocharger operating point by varying speed and load conditions. The air flow through the  
236 compressor was measured directly using an *ABB Sensyflow* hot wire flow meter. The flow through  
237 the turbine was estimated from the compressor air flow and the fuel flow, measured from a *CP*  
238 *Engineering FMS1000* gravimetric fuel balance. Pressure measurements were made on a 10Hz basis  
239 using *Druck PTX* sensors at the inlet and outlet of both turbine and compressor. At the inlet and  
240 outlet of the turbine pressures were also measured on an engine crank angle basis using *Kistler 4049*  
241 piezo-resistive sensors.

242 The following experiments were repeated three times to increase confidence in results:

- 243 • 14 thermally stable conditions following 8min stabilisation period (figure 6a)

244 • Transient step experiments (series of step changes in engine operating point with three  
 245 minute hold time (see figure 6b). Three minutes was chosen as this allows for the system to  
 246 stabilise between transients thus allowing for the individual analysis of each step transient.

247 The steady and transient engine operating conditions are also shown on the turbine and compressor  
 248 maps in figure 5.

#### 249 4.4 Measurement Uncertainty

250 Measurement uncertainty has been carried out for total heat transfer and convection coefficients.  
 251 This combines uncertainty of individual measurement instruments (see table 2) into the calculated  
 252 values used in the results section using equation 18 [28].

$$u_y = \left( \sum_{i=1}^n \left( \frac{\partial y}{\partial x_i} u_{x_i} \right)^2 \right)^{0.5} \quad 18$$

253

254 As an example, applying equation 18 to equation 4 for the uncertainty of the heat transferred from  
 255 gas to turbine casing yields:

$$u_{Q_{G/T}} = \left( \begin{array}{l} (\dot{m}_t c_{p,T} u_{T_{03}})^2 + (-\dot{m}_t c_{p,T} u_{T_4})^2 \\ + (c_{p,t} (T_{03} - T_{04}) u_{\dot{m}_t})^2 \\ + (\dot{m}_t (T_{03} - T_{04}) u_{c_{p,t}})^2 \\ + (-u_{W_t})^2 \end{array} \right)^{0.5} \quad 19$$

256 Where the uncertainty in the work transfer,  $u_{W_t}$ , is estimated in a similar manner by applying  
 257 equation 18 to equations 5-7. Uncertainties arising solely as a result of sensor uncertainties are  
 258 given as the solid square points in the upper graphs of figure 7 (a-c) for turbine work, turbine heat  
 259 transfer rate and turbine Nusselt number. The uncertainty for turbine work increases at lower  
 260 turbine speeds because the measurement is dependent on the difference in temperatures before  
 261 and after the compressor: as the speed and power reduce, this difference becomes considerably  
 262 small. In contrast, heat transfer rate and Nusselt number have increasing uncertainties at higher

263 shaft speeds. This is because the uncertainties for both these quantities is highly sensitive to mass  
 264 flow rate and therefore higher uncertainty results at higher mass flows.

265 **Table 2: Measurement accuracy of various quantities measured on the gas stand/engine stand**

Measurement	Unit	Sensor	Accuracy
Temperature	°C	k-type thermocouple	+/-2°C
Air Mass flow	kg/h	ABB Sensyflow	<1%
Pressure	kPa	Piezo-resistive	+/-0.04%

266

267 In addition to the measurement uncertainty, of particular importance is the calculation of  
 268 compressor work based on the temperature measurements at compressor inlet and outlet (equation  
 269 6). As stated, this approach ignores heat transfer effects in the compressor for the determination of  
 270 heat transfer in the turbine. To quantify the uncertainty in the proposed approach, it will be  
 271 assumed that the magnitude of heat transfer rate in the compressor is similar to that presented by  
 272 Serrano et al [16]: the ratio of heat transfer rate to total enthalpy rate change in the compressor was  
 273 presented as a function of total enthalpy rate change in the turbine (equation 20). The estimated  
 274 compressor heat transfer is directly equated to an additional uncertainty source for compressor  
 275 work equation 21).

$$\frac{\dot{Q}_c}{\Delta\dot{H}_c} = f(\Delta\dot{H}_t) \quad 20$$

$$u_{w_c,HT} = \dot{Q}_c \quad 21$$

276

277 The uncertainty due to heat transfer is combined with the sensor uncertainties to calculate the  
 278 influence on key uncertainties using equation 18. These are presented alongside the sensor only  
 279 uncertainties in figure 7. For turbine work, the increased uncertainty as a result of ignoring heat  
 280 transfer is significant for turbocharger speeds below 100krpm and at 50krpm the increased  
 281 uncertainty through ignoring heat transfer is around 30% of the measured value. Clearly this is a

282 severe limitation and illustrates why turbine maps measured in this way are only provided at higher  
283 shaft speeds. The effect on uncertainty for turbine heat transfer and Nusselt number is considerably  
284 less and at 50krpm the uncertainty increase is only 6% of the measurement. It is these latter two  
285 quantities that are most important and high uncertainties for the turbine work will be tolerated.

## 286 **5 RESULTS AND DISCUSSION**

### 287 **5.1 Overview of Heat transfer**

288 The ratio of heat to work transfer gives an indication of the importance of heat transfer for turbine  
289 performance prediction. This is shown over the engine operating map in figure 8, against turbine  
290 mechanical power in figure 9 and over the turbine map in figure 10. This highlights the problem that  
291 heat transfer is more significant a lower turbine powers where they are not typically mapped,  
292 corresponding to lower engine powers. The results in figures 9 and 10 are obvious because they  
293 show that heat transfer is strongly linked to temperature and operating point. It is interesting to  
294 note that heat transfer accounts for at least 20% of the enthalpy drop over the turbine but that at  
295 low turbine powers, this proportion can be significantly higher, even with low turbine gas  
296 temperatures. Through figure 9, the exponential correction curves used to correct turbine maps for  
297 correlating engine models to measured data [8, 11] are clearly visible.

### 298 **5.2 Internal Convection**

#### 299 **5.2.1 Steady Flow Results**

300 Measured Reynolds and Nusselt numbers are plotted for different VGT positions and turbine inlet  
301 temperatures in figure 11; 95% confidence intervals are also shown. The results show that the range  
302 of Re numbers under cold and hot flow conditions are an order of magnitude different due to  
303 changes in density and mass flow. For a turbine inlet temperature of 500°C (figure 11 a), the Nu/Re  
304 relationship has a similar shape to the Seider-Tate correlation for straight pipes. There is also very  
305 little distinction within the error margins with respect to VGT position. The 95% uncertainty margins



306 are acceptable but grow with increasing Reynolds number. In contrast, at 100°C TIT, (figure 11 b),  
307 the measured Nu/Re relationship has an exponential shape and very high uncertainty. This is  
308 because the largest uncertainty is associated with the temperature measurements, and the  
309 sensitivity of this uncertainty is linked to mass flow (equation 8). As mass flow is higher under colder  
310 conditions, the uncertainty is also higher.

311 The data measured for turbine inlet temperature of 500°C was used to fit coefficients to equation  
312 12. Initially these correlations have been established using the gas stand specific temperature  
313 measurements rather than the thermocouples mounted onto the turbocharger as this provided  
314 consistency with other published works, assumed to use these measurements, which allows a direct  
315 comparison. In figure 12, these new convective correlations are compared to the other published  
316 correlations previously presented in table 1. Cormerais et al. [18] and Reyes [22] established  
317 turbocharger-specific correlations (three different devices for Reyes) resulting in an individual  
318 correlation for each device. These devices were of varying size but all aimed at passenger car  
319 applications ranging from 1.2 to 2.0L displacement. Baines et al [7] derived a single correlation,  
320 validated for three turbochargers of similar size for automotive truck applications (hence larger than  
321 those studied by Reyes and Cormerais et al.). In contrast, Romagnoli et al. [19] proposed to use an  
322 established correlation for flow in pipes.

323 It can be seen that the pipe flow correlation (Romagnoli) agrees well with two of the correlations  
324 proposed by Reyes (1 and 2). In contrast the Reyes 3 correlation estimates higher convective heat  
325 transfer and agrees better with the correlation proposed by Cormerais et al. Baines' correlation sits  
326 in between these two extremes. The correlations derived from the present work give a similar  
327 magnitude to the correlations from Cormerais and Reyes (3). There is significant variation  
328 depending on the VGT guide vane angle which could be expected because of the way these guide  
329 vanes will affect the flow characteristics within the volute and wheel. It should be remembered that  
330 the coefficients of the correlation are required to capture the complex 3D flow phenomenon

331 occurring within the device. Reyes [22] also observed this influence of VGT with similar magnitudes  
332 difference. However, in this study it is observed a gradual decrease with VGT opening, Reyes  
333 observed a correlation with peak turbine efficiency, meaning peak convective heat transfer was  
334 observed at intermediate VGT vane angles with lower correlations as the vanes were opened or  
335 closed.

336 There are no obvious links between the correlations and the characteristics of each of the  
337 turbochargers based on the available data.

338 1. When considering the size of the different turbochargers, the correlations proposed by  
339 Reyes [22] rank with device size, with Reyes 1 the largest (from 2.0L engine) and Reyes 3  
340 smallest from a 1.2L engine. However the correlation proposed by Baines [7] is for  
341 considerably larger commercial vehicles and the device used in the present paper is taken  
342 from a 2.2L engine.

343 2. When considering the inclusion of variable geometry guide vanes, the turbochargers used to  
344 derive correlations Reyes 1 and 2, Cormerais and those from the present paper were all  
345 fitted with the VGT devices and span the full range of observed convective heat transfer.

346 The wide range of values proposed in the literature show that it is not yet possible to derive a single,  
347 simple convective heat transfer correlation applicable across different turbocharger devices. Whilst  
348 this highlights the need for further study in this area, for the purpose of the present study it  
349 demonstrates consistency with other research findings.

350 Turbine Nusselt numbers were also calculated based on the gas stand data using the thermocouples  
351 located on the turbocharger inlet and outlet port, and identical to the measurements used on-  
352 engine. These are compared to those obtained previously and to results from the engine  
353 experiments in figure 13. These will be discussed in the next section.

354 **5.2.2 On-engine results**

355 The analysis was repeated for engine based experiments and results are shown in figure 13. To allow  
356 direct comparison of engine conditions with exhaust gases and gas stand conditions with air, the  
357 results have been adjusted to a Prandtl number of 0.7 using equation 22.

$$Nu_{(Pr=0.7)} = Nu_{(Pr=x)} \frac{0.7^{1/3}}{Pr_x^{1/3}} \quad 22$$

358

359 The results in figure 13 shows good agreement between the engine and gas stand data when the  
360 same temperature sensors are used in both cases. The agreement is worst when comparing the  
361 engine data with that obtained from the gas stand using gas stand standard instrumentation,  
362 especially at low Reynolds numbers. The discrepancy between gas stand and engine Nusselt number  
363 at low Reynolds number can primarily be attributed to the measurements of temperatures at inlet  
364 and outlet of the turbine. On the gas stand, conditions for this measurement are more favourable as  
365 the flow is both steady and enters and leaves the turbine long straight pipes. On the engine, the  
366 flows are both pulsating and the geometry of the exhaust manifold and subsequent exhaust system  
367 make the temperature measurement particularly complex. Future studies should consider further  
368 measures to promote the accuracy in an engine situation such as the inclusion of bespoke  
369 measurement sections replicating to some extent the gas stand layout. It is interesting to note that  
370 when using the same instrumentation on-engine and on gas stand, there is strong agreement  
371 despite the differences in flow conditions showing that sensor location is a dominant effect.

372 Using the gas-stand derived correlation with standard instrumentation to predict the turbine outlet  
373 temperature on-engine, improved turbine outlet temperature prediction as shown in figure 14. This  
374 equates to a reduction of mean error from 33°C to -3°C but an increase in standard deviation of  
375 error from 10°C to 20°C.

376 The model was also used to investigate the effect of pulsating flows on the internal convection. Until  
377 now, the Re and Nu numbers have been considered both spatially and time averaged. Pulsating  
378 flows equate to pulsating Reynolds numbers as illustrated in figure 15 (obtained from the 1D turbine  
379 flow model). Figure 16 compared three different Nusselt/Reynolds relationships obtained as follows:

- 380 1. Raw Measurement: obtained using the raw measurements of mass flow from the engine  
381 test bench (from the intake air flow and fuel flow measurements).
- 382 2. Pulsating Simulation: obtained by using a simulated mass flow based obtained by applying  
383 measured instantaneous pressure measurements from the turbine inlet and outlet and  
384 applying these to a double orifice model proposed by Serrano et al. [26].
- 385 3. Arithmetic mean of pulsating simulation: average Nusselt and Reynolds number over two  
386 engine revolutions corresponding to four pulsations, one from each of the engine cylinders.

387 There is a small offset between the arithmetic mean and the point that results from the physical  
388 averaging due to slow sensors response. However, as shown in figure 14, this has little effect on the  
389 accuracy of turbine outlet temperature prediction.

### 390 **5.3 Effect of transient operation**

391 The transient events on engine affect the turbocharger at a range of timescales [29]. It has been  
392 shown in the previous section that the flow pulsations have only a small impact on the heat transfer  
393 phenomenon. This is explained by the large time constant associated with the thermal inertia of the  
394 turbine housing compared with the frequency of the exhaust pulsations. Engine transients occur  
395 over a longer timescale and the dynamics of heat transfer will become significant as illustrated in  
396 figure 17. The figure shows measured turbine inlet temperature and rotational speed and calculated  
397 heat flows following two near step changes in engine power. The calculated heat flows are obtained  
398 by the heat transfer model whilst applying measured boundary conditions of turbine inlet and outlet  
399 pressure, turbine inlet temperature and turbine rotational speed.

400 Figure 17 (a) illustrates the case of a step up in engine power. Initially the engine is operating under  
401 stable conditions: heat flow from gas to the turbine housing almost equals the heat flow out of the  
402 housing and the net heat flow into or out of the housing wall (heat storage) is null. The engine  
403 undergoes a rapid power transient at time  $t=0s$  causing a rise in turbine inlet temperature through  
404 higher fuelling levels. Other control actions also occur such as closing of variable geometry guide  
405 vanes. The turbine rotation speed increases sharply over the next 5 seconds with the dynamics  
406 controlled by the mechanical inertia of the shaft and loadings on the compressor. The measured  
407 turbine inlet temperature undergoes two distinct phases: over the first 30s the temperature rise is  
408 relatively large, increasing from  $320^{\circ}C$  to  $474^{\circ}C$ . The temperature overshoot in this period can be  
409 explained by a particular transient characteristic of the engine control which allows a temporary  
410 overboosting of the engine. From 30s to 150s after the change in engine power, the temperature  
411 rise is much slower with an increase to  $492^{\circ}C$ . These temperature measurements, which are also  
412 used to calculate the heat flows in figure 17 will be affected by the response of the thermocouples  
413 used in the measurement. The response of the thermocouple is dependent on the thermal inertia,  
414 and therefore the size of the thermocouples, but also the flow conditions and notably the Reynolds  
415 number in the pipe. Analysis presented in appendix 1 for 3mm diameter thermocouples shows that  
416 the 95% response time for turbocharger inlet and outlet conditions is in the range of 4-20s. Clearly  
417 this will affect the accuracy during the initial rapid response corresponding to the first 30s of the  
418 response. However, this will have a much smaller impact on the accuracy of the second phase of the  
419 response 30-150s. The heat flow from gas to housing peaks at the beginning of the transient (in this  
420 case at around 7kW) before slowly falling to a value of around 3.6kW three minutes later. This spike  
421 in heat flow is accounted for by the accumulation of heat in the turbine housing as it warms up.

422 Figure 17 (b) shows the opposite case for a step down in engine load. The heat loss from the gas to  
423 housing has a minimum which, in this particular case, is -200W before tending to around +200W  
424 three minutes after the transient. This indicates that the heat flow is reversed just after the

425 transient, flowing from the housing to the gas. The same comments for figure 17 (a) with regard to  
426 thermocouple response time equally apply to figure 17 (b).

427 For a series of steps in engine power of different magnitude, the peak and settled heat transfers  
428 have been defined as follows:

- 429 •  $\Delta T_3$ : The temperature difference between instances just prior to step change in engine  
430 power and three minutes after the step change.
- 431 •  $\dot{Q}_{Peak}$ : Maximum or minimum measured heat flow during first 60s following step change in  
432 engine power
- 433 •  $\dot{Q}_{Settle}$ : Mean measured heat flow between 160 and 180 seconds following step change in  
434 engine power

435 The quotient of these two heat flow values have been plotted against step change in turbine inlet  
436 temperature in figure 18.

437 For heat flow between the exhaust gases and the turbine casing, figure 18a shows that larger TIT  
438 transients result in larger ratio of initial to settled heat transfer. It is important to bear in mind that  
439  $Q_{settle}$  was always a positive, as ultimately the exhaust gases are always hotter than the ambient.

- 440 • Points in the upper right quadrant correspond to a positive step in exhaust gas temperature.  
441 Every case in this quadrant has a  $Q_{peak}/Q_{settle}$  ratio greater than 1 meaning the peak heat flow  
442 is larger than the settled flow.
- 443 • Points located in the upper left quadrant correspond to step reductions in TIT temperature  
444 and all have a  $Q_{peak}/Q_{settle}$  value between 0 and 1. These points all correspond to step  
445 reductions in exhaust gas temperature of less than 100°C.
- 446 • Points located in the lower left quadrant correspond to situations where the peak heat flow  
447 is reversed (i.e. from the casing to the exhaust gases). These correspond to larger reductions  
448 in TIT of more than 100°C.

449 Figure 18b shows the same ratio for heat transfer from the turbine housing to its surroundings  
450 (ambient and bearing housing). For all steps, positive and negative, of less than 100°C this ratio is  
451 approximately equal to 1 meaning there is no significant peak in heat flow. This is because the  
452 turbine housing temperature does not change significantly and ambient temperature remains  
453 constant, therefore heat transfers with the surroundings are not affected. It is only when much  
454 larger step changes in turbine inlet temperature are induced that peak heat flow begins to appear.

## 455 **6 CONCLUSIONS**

456 In this paper, an experimental and lumped capacity modelling approach was used to assess heat  
457 transfer characteristics in turbocharger turbines. Through this work, the following conclusions have  
458 been drawn:

- 459 1. Heat transfer in the turbine always represents at least 20% of enthalpy change in the  
460 turbine, however it can be significantly more under low turbine power conditions. This  
461 corresponds to the low power operating conditions of the engine in the low speed/low  
462 torque region
- 463 2. It is difficult to compare the fitted curves for Nusselt /Reynolds correlations at different  
464 turbine inlet temperatures because the Reynolds numbers vary by an order of magnitude  
465 due to changes in fluid density. Consequently direct comparison relies on considerable  
466 extrapolation away from the measured data.
- 467 3. Heat transfer correlations determined on-engine and gas stand can be significantly different  
468 due to different instrumentation layouts, however when consistent sensors are used across  
469 facilities, good agreement is obtained despite the differences in flow conditions.
- 470 4. Despite the discrepancies incurred due to instrumentation differences, the use of heat  
471 transfer correlations obtained from the gas stand to simulate on-engine conditions will  
472 provide a significant improvement in prediction accuracy using either averaged or pulsating  
473 flow Reynolds numbers. This shows quasi-steady behaviour can be assumed for convective

474 heat transfer coefficient. Care should be taken to account for the change in Prandtl number  
 475 due to variations in gas composition between air and exhaust gases. Application of the gas  
 476 stand derived heat transfer correlation reduced the turbine outlet temperature prediction  
 477 error from 33°C to -3°C.

478 5. Operation under transient conditions shows that the thermal inertia of the housing  
 479 significantly influences the heat flow because of the change in temperature difference  
 480 between gas and wall. For large reductions in turbine inlet temperature (greater than  
 481 100°C), the heat flow was reversed during the transient.

## 482 7 GLOSSARY

### 483 7.1 Nomenclature

A	Area	$m^2$
c	Empirical Constant	
$C_{fr}$	Friction Constant	$W/rpm^2$
$c_p$	Heat capacity at constant pressure	$J/kgK$
D	Diameter	m
h	Specific enthalpy, Convective Heat Transfer coefficient	$kJ/kg,$ $W/m^2K$
k	Thermal Conductivity	$W/mK$
L	Characteristic length	m
m	Mass	kg
$\dot{m}$	Mass flow	kg/s
N	Shaft Speed	1/min
Nu	Nusselt Number	
P	Pressure	Bar
Pr	Prandtl Number	



$q$	Specific heat flow	$\text{kJ/kg}$
$\dot{Q}$	Heat transfer rate	$\text{W}$
$Re$	Reynolds Number	
$T$	Temperature	$\text{K}$
$t$	time	$\text{s}$
$u$	Uncertainty	
$v$	velocity	$\text{m/s}$
$W$	Work	$\text{J}$
$\dot{W}$	Work transfer rate	$\text{W}$
$x$	Length (Conduction)	$\text{m}$
	Independent Variable (uncertainty analysis)	
$y$	Dependent Variable (uncertainty analysis)	
$\gamma$	ratio of specific heats	
$\eta$	Efficiency	
$\mu$	Dynamic Viscosity	$\text{kg s}^{-1} \text{m}^{-1}$
$\rho$	Density	$\text{Kg/m}^3$
$\tau$	Thermocouple Time Constant	$\text{s}$

484

## 485 7.2 Subscripts

0	Stagnation (temperature)
1	Pre Compressor
1'	Pre compression
2	Post Compressor
2'	Post Compression
3	Pre Turbine
3'	Pre Expansion

4	Post Turbine
4'	Post Expansion
a	After compression/expansion
act	Actual
b	Before compression/expansion
B	Bearing Housing
bulk	Fluid bulk property
c	Compressor
conv	Convective
corr	Corrected
d	inlet diameter
f	friction
g	Gas
rad	Radiation
s	Isentropic (Efficiency)
skin	Fluid film skin property
t	Turbine
tc	thermocouple
T	Turbine Housing
work	Mechanical Work

486

### 487 7.3 Abbreviations

BMEP	Brake Mean Effective Pressure
CO <sub>2</sub>	Carbon Dioxide
EGR	Exhaust Gas Recirculation
RMSE	Root mean square error
TIT	Turbine Inlet Temperature
VGT	Variable Geometry Turbine

488

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578 **9 APPENDIX. ESTIMATING THE TRANSIENT RESPONSE OF THERMOCOUPLE**  
579 **PROBES**

580 In order to estimate the transient response of a thermocouple probe situated in the gas stream at  
581 the inlet or outlet of the turbocharger turbine, the problem can be assumed to be equivalent to that  
582 of a small sphere located within a flow of gas. The convective heat transfer coefficient for such a  
583 configuration has been determined empirically and can be calculated from equation 23 [21].

$$Nu_{tc} = 2 + (0.4\sqrt{Re_{tc}} + 0.06Re_{tc}^{2/3})Pr^{0.4} \left( \frac{\mu_{bulk}}{\mu_{skin}} \right)^{1/4} \quad 23$$

584 Where the characteristic length is the diameter of the thermocouple tip  $d_{tc}$ .

585 The main stream velocity at the inlet or outlet of the turbine can be related to the mass flow in the  
586 turbine using equation 24. This velocity is then used in the calculation of the Reynolds Number.

$$v = \frac{\dot{m}_t}{\rho_g A} \quad 24$$

587  
588 The heat transfer coefficient can then be derived from the Nusselt number and the fluid properties  
589 and be used to estimate the heat transfer by convection between the gas and the thermocouple tip  
590 using Newton's law of cooling (equation 25).

$$\dot{Q}_{g,tc} = h_{tc} A_{tc} (T_g - T_{tc}) \quad 25$$

591  
592 Assuming that heat transfer by convection between the gas and the thermocouple is the only  
593 significant heat flow, then this heat transfer can be equated to the temperature rise of the  
594 thermocouple tip using equation 26

$$m_{tc} c_{p,tc} dT_{tc} = h_{tc} A_{tc} (T_g - T_{tc}) dt \quad 26$$

595

596 Equation 26 can be rearranged into equation 27 which defined a one degree of freedom first order  
597 system.

$$\frac{m_{tc}c_{p,tc}}{h_{tc}A_{tc}} \frac{dT_{tc}}{dt} + T_{tc} = T_g \quad 27$$

598

599 Considering that the system is initially at equilibrium where  $T_{tc} = T_g(0)$ . At time  $t=0$ , a step change  
600 in gas temperature  $\Delta T_g$  occurs. The thermocouple response would be given by equation 28.

$$T_{tc}(t) = T_g(0) + \Delta T_g \left(1 - e^{-\frac{t}{\tau}}\right) \quad 28$$

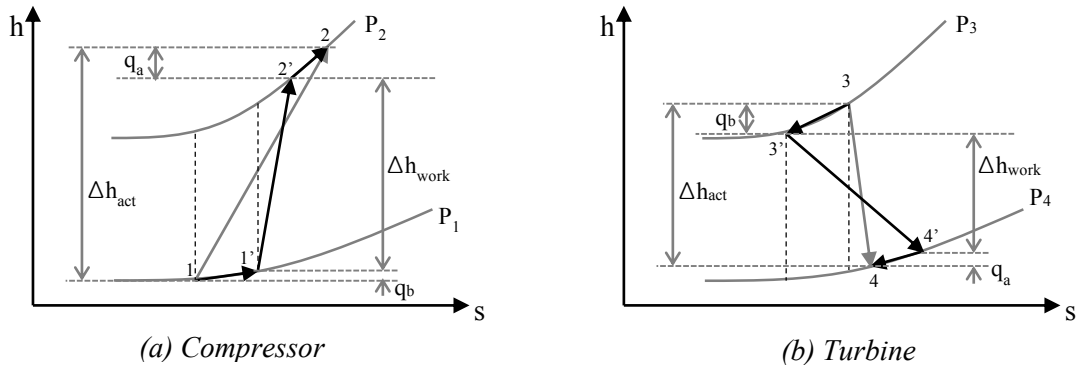
601 Where the time constant  $\tau = \frac{m_{tc}c_{p,tc}}{h_{tc}A_{tc}}$

602 This standard first order system response has a 95% response time of three time constants. This  
603 response time is shown for thermocouples of diameter 0.5mm, 1.5mm and 3mm as a function of  
604 turbine Reynolds number in figure 19. These have been calculated assuming that the fluid is air at  
605 400°C. It should be noted that this Reynolds number ( $Re_T$ ) is related to the turbine inlet or outlet  
606 diameter and therefore directly comparable to the heat transfer correlations used in this work. This  
607 is different to the Reynolds number used in this appendix ( $Re_{tc}$ ) which is used to determine the heat  
608 transfer coefficient to the thermocouple tip. The results show that for a 3mm diameter  
609 thermocouple, the 95% response time will vary between 5s at high flow rates to 17s at low flow  
610 rates. Variations as a result in the change in fluid properties due to changes in temperature affect  
611 these results by around 8% per 100°C. This means that for the same 3mm thermocouple, the range  
612 of 95% response times at the highest likely turbine inlet temperature of 900°C is 4s to 14s whilst at  
613 the lowest likely turbine inlet temperature of 200°C it is 6s to 20s.

614 It can be concluded from this analysis that the response time of the thermocouples used in this work  
615 will be in the region of 4 to 20s.

616





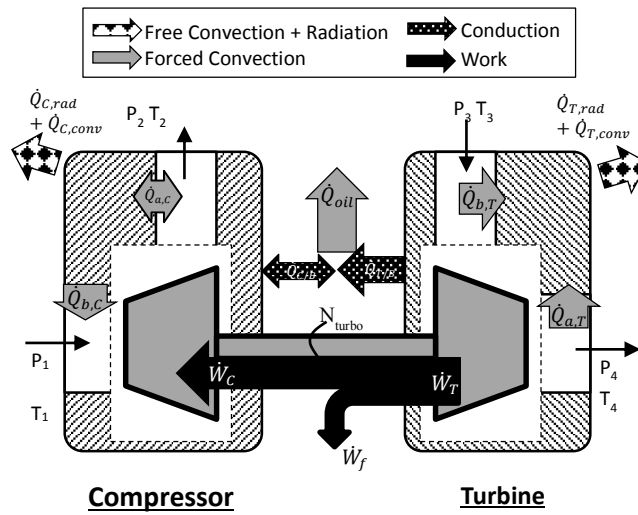
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Figure 1: Apparent and assumed compression and expansion processes in (a) compressor and (b) turbine

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Figure 2: Overview of turbocharger heat transfer model

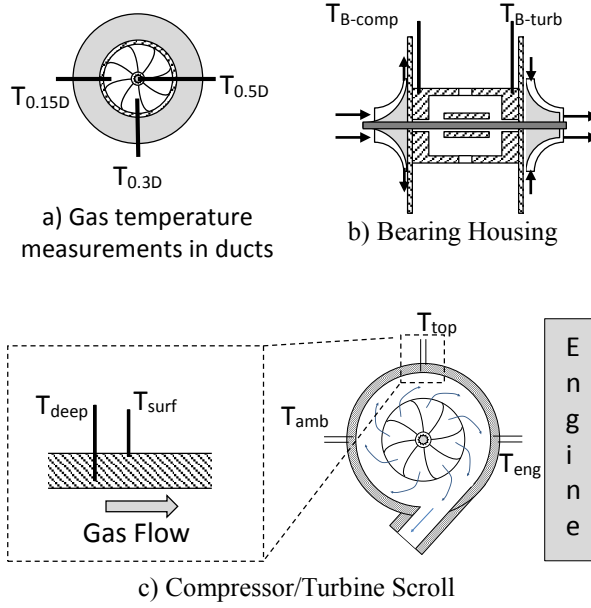


Figure 3: Thermal instrumentation of turbocharger housing and gas ports

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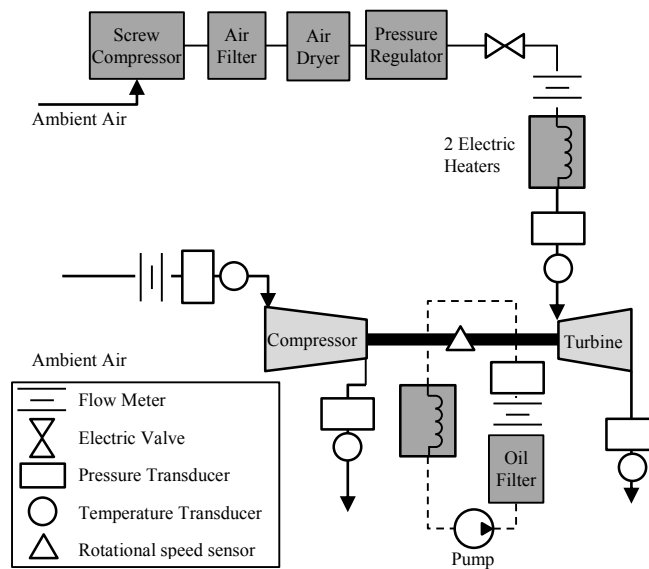
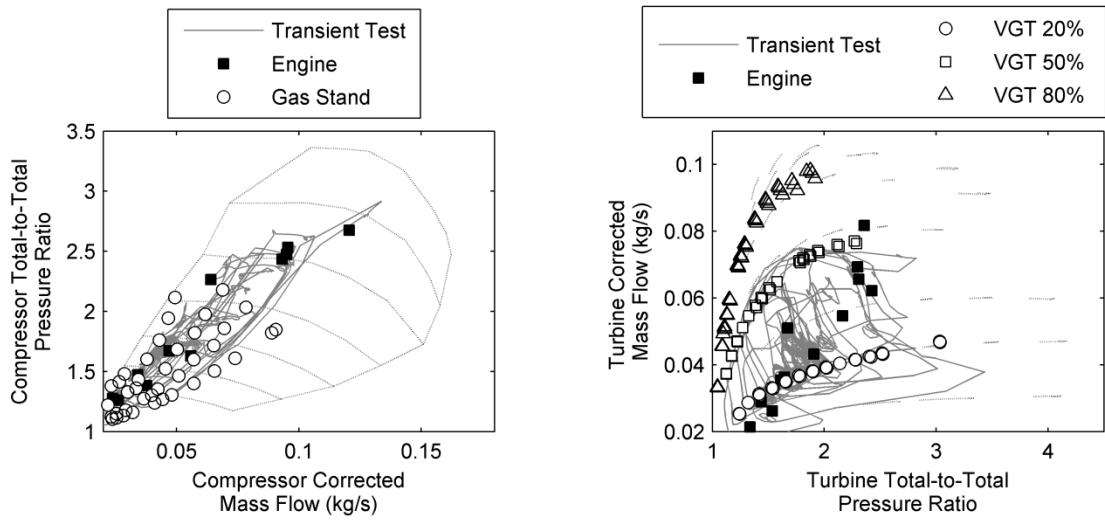


Figure 4: Schematic of gas stand installation

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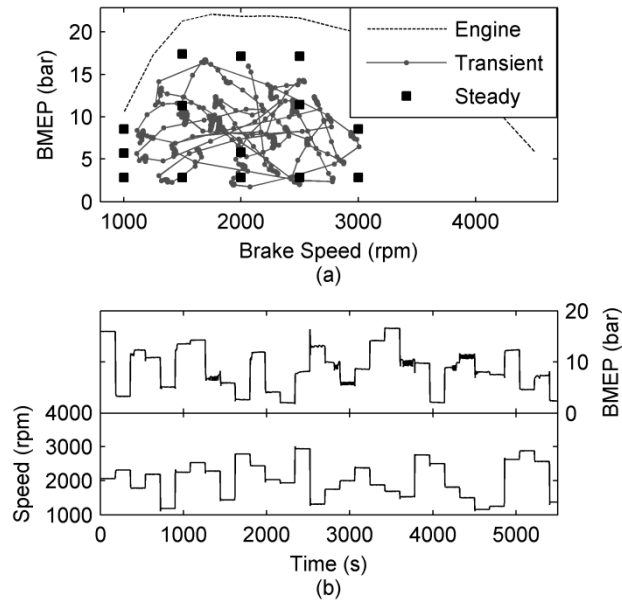
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Figure 5: Compressor and Turbine operating points during steady flow gas stand experiments

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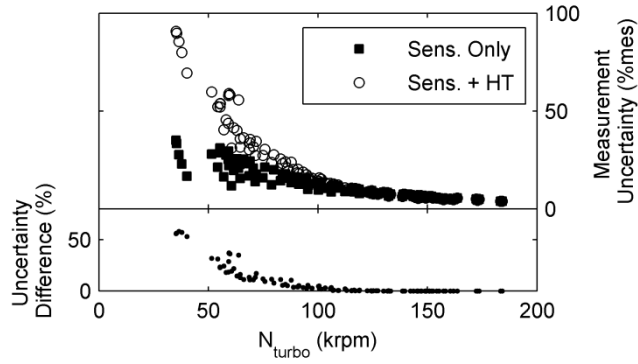


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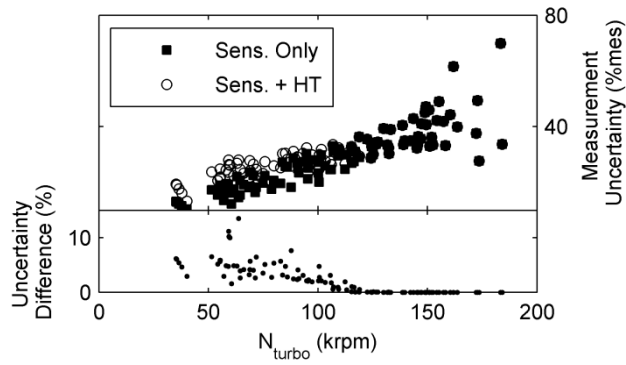
Figure 6: Compressor, Turbine and engine operating points for steady and transient experiments

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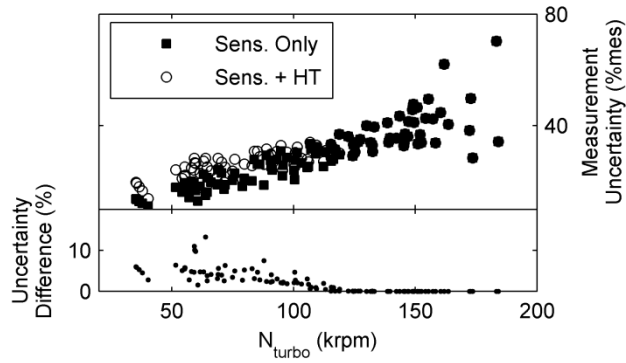
(a) - Turbine Work

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(b) - Turbine Heat Transfer

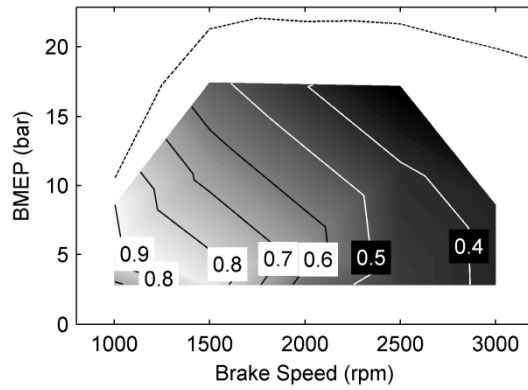
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(c) - Turbine Nusselt Number

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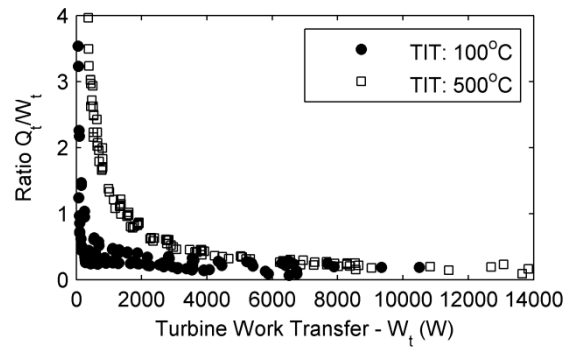
641 **Figure 7: Estimated combined uncertainty for (a) turbine work transfer rate, (b) turbine heat transfer rate and (c)**  
 642 **turbine Nusselt number. Uncertainties due only to sensor uncertainty are compared with uncertainty due to sensors**  
 643 **and ignoring heat transfer in the compressor side.**



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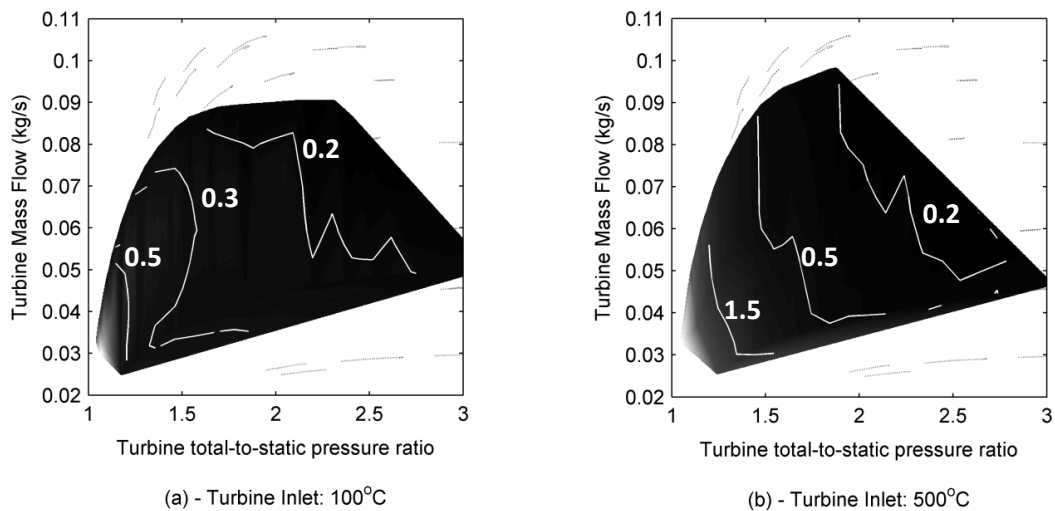
Figure 8:  $Q_{G/T}/W_T$  over the engine operating map



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Figure 9: Ratio of turbine total heat loss to turbine work from gas stand testing with TIT 100°C and 500°C

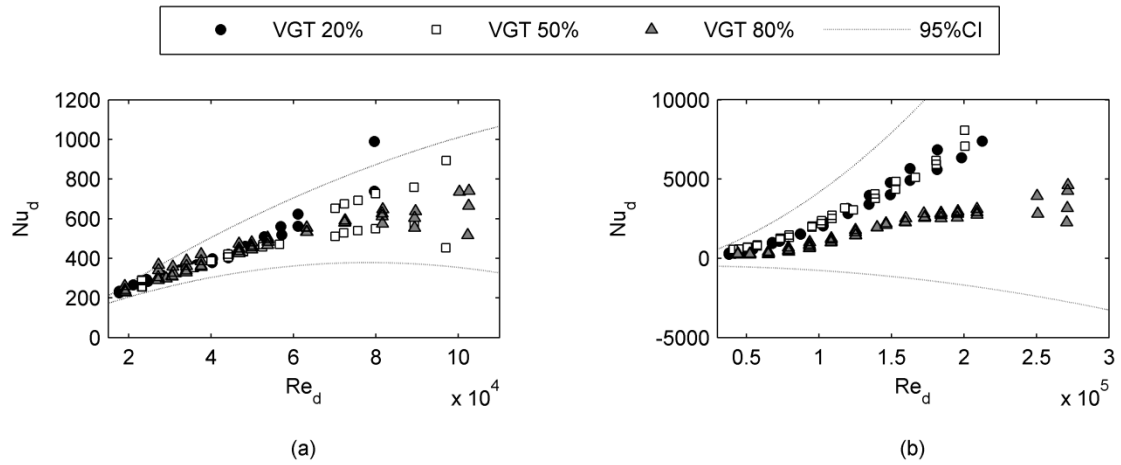


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Figure 10: Ratio of turbine total heat loss to turbine work over the turbine operating map for turbine inlet temperature (a) 100°C and (b) 500°C

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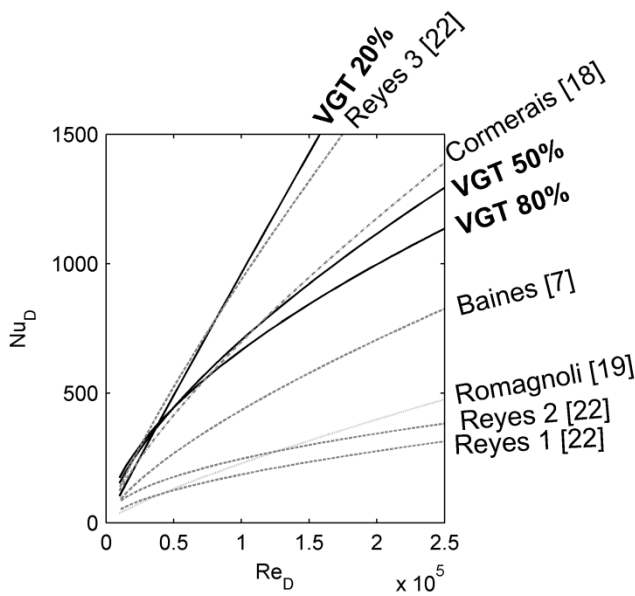


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653 **Figure 11: Spatially averaged Nu/Re correlation based on steady flow gas stand experiments for turbine inlet**  
 654 **temperature (a) 500°C and (b) 100°C**

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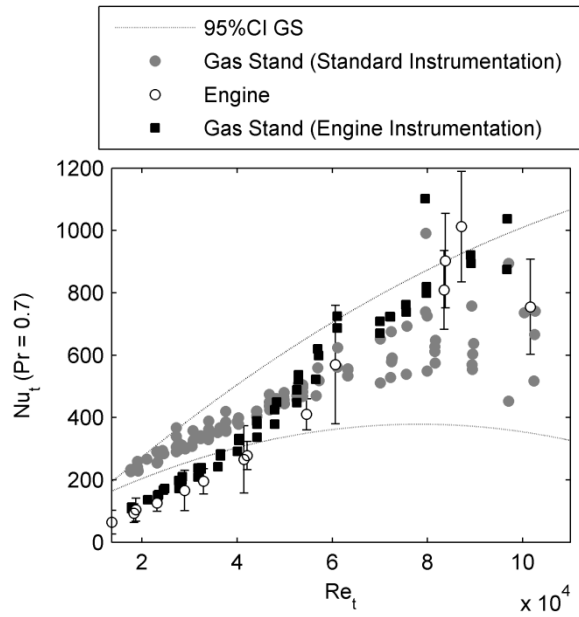
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658 **Figure 12: Comparison of fitted convection correlations with selected correlations from published literature**

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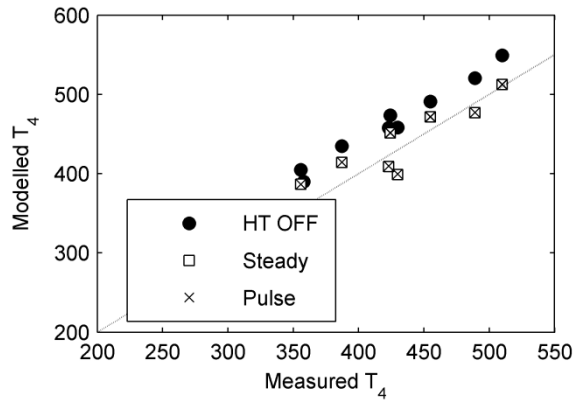


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661 **Figure 13: Spatially-Time averaged Nu/Re correlation for engine experiments assuming air and exhaust gases,**  
 662 **corrected for Pr=0.7**

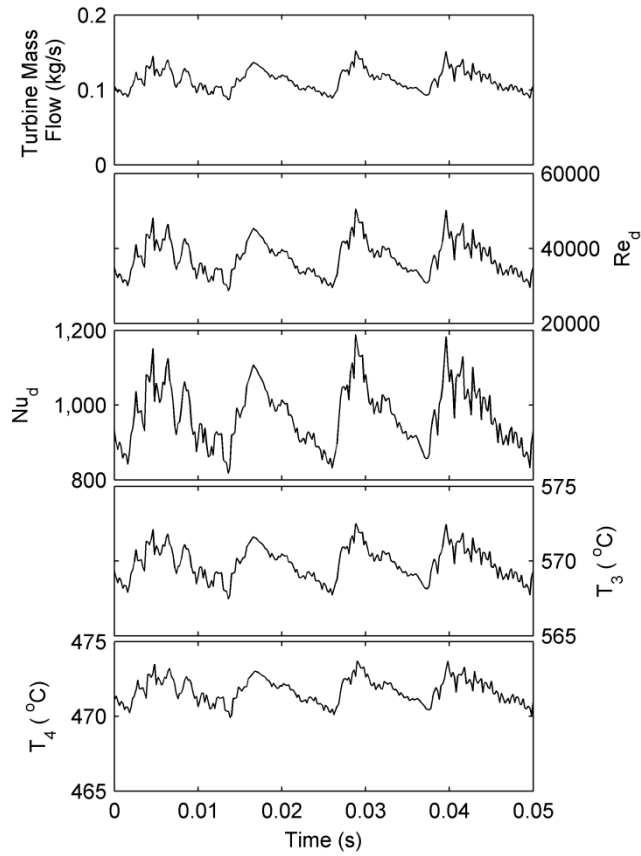
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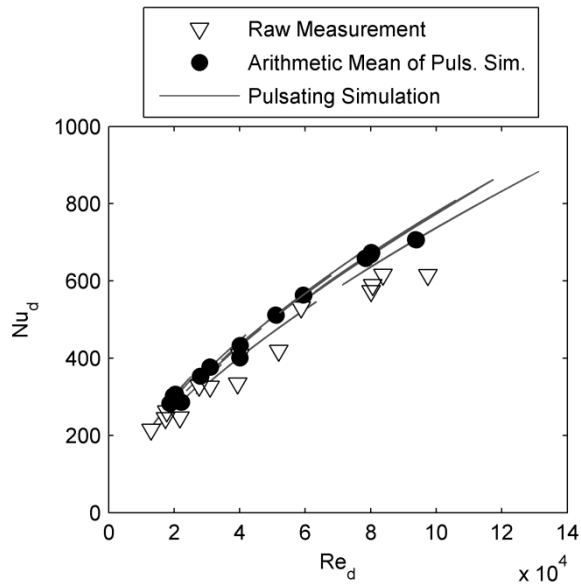
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666 **Figure 14: Prediction of turbine outlet temperature without heat transfer, using Spatially-Time-averaged Reynolds**  
 667 **number and using Spatially-averaged instantaneous Reynolds number**



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669 **Figure 15: Predicted turbine flow rate, constant gas temperature, instantaneous Reynolds and Nusselt numbers and**  
 670 **predicted pulsating turbine outlet temperature**

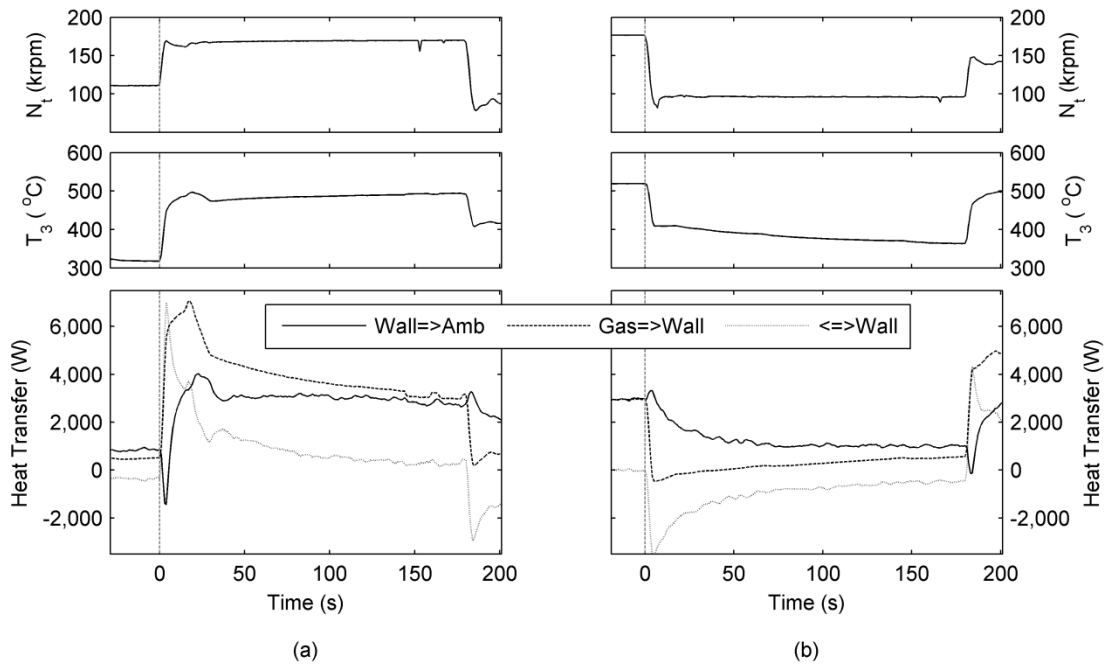


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672 **Figure 16: Nusselt/Reynolds relationship for Spatially-Time averaged Reynolds number and spatially averaged**  
 673 **instantaneous Reynolds number**



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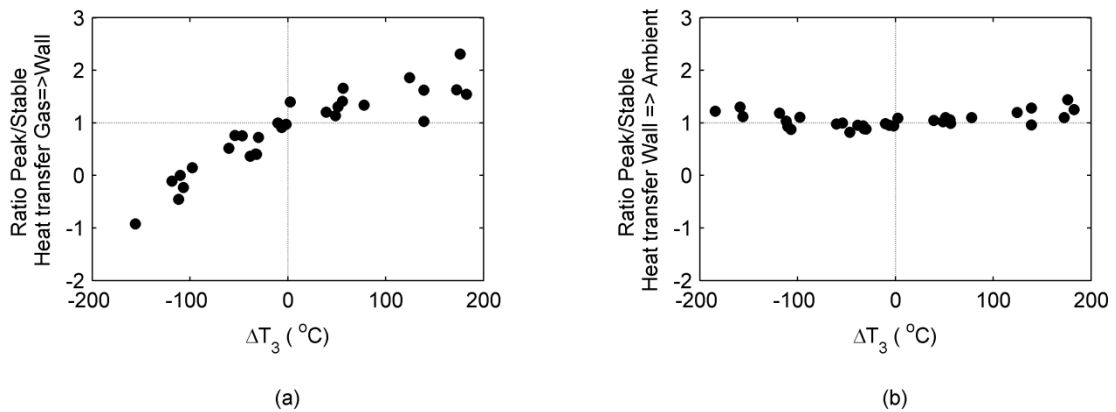


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676 **Figure 17: Evolution of turbine rotor speed, Turbine inlet temperature and heat flows between wall and ambient,**  
677 **exhaust gas and wall and heat storage in thermal capacitance of housing during (a) step up in engine power and (b)**  
678 **step down in engine power (0 on the x axis indicates the beginning of each transient)**

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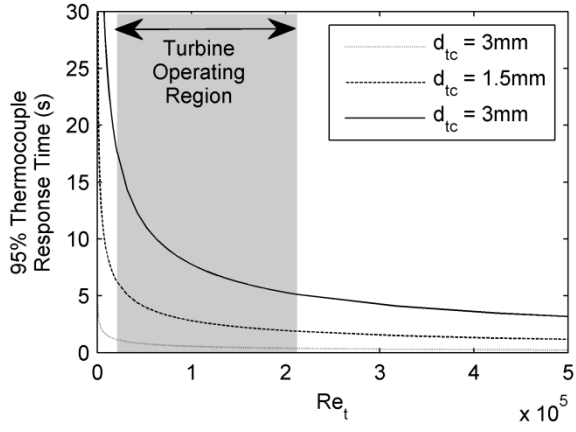
682 **Figure 18: Quotient of Peak transient heat flow to settled heat flow between (a) exhaust gas and turbine housing and**  
683 **(b) turbine housing and ambient**

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689 **Figure 19: Estimated Thermocouple 95% response time for thermocouple diameters 0.5mm to 3mm as a function of**  
690 **turbine Reynolds number**

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